

Research Article

Order and Chaos in a Prey-Predator Model Incorporating Refuge, Disease, and Harvesting

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In this paper, a mathematical model consisting of a prey-predator system incorporating infectious disease in the prey has been proposed and analyzed. It is assumed that the predator preys upon the nonrefugees prey only according to the modified Holling type-II functional response. There is a harvesting process from the predator. The existence and uniqueness of the solution in addition to their bounded are discussed. The stability analysis of the model around all possible equilibrium points is investigated. The persistence conditions of the system are established. Local bifurcation analysis in view of the Sotomayor theorem is carried out. Numerical simulation has been applied to investigate the global dynamics and specify the effect of varying the parameters. It is observed that the system has a chaotic dynamics.

1. Introduction

The interaction between a predator and their prey is one of the most important interaction types in ecology as well as applied mathematics due to its wide existence in our real life. Although such interaction models may appear to be simple mathematically, they are very challenging and complicated [1]. Since the classical Lotka (1925) and Volterra (1926) model on prey-predator interaction, which was developed on the basis of chemical principle of mass action, many research works have been done in literature. Later, the simple growth and decay terms in the Lotka-Volterra model were replaced by more complex functions such as the logistic growth; prey-dependent functional responses such as Holling's I, II, and III types; ratio-dependent functional response, and predator-dependent type of functional response, see, for example, [2-8] and the references therein. It is well known that the two species models are very restrictive in their dynamics and can give rise to only two basic dynamics approaches to equilibrium or to a limit cycle, while ecological situations demonstrate very complex dynamics in nature. Accordingly, multispecies systems comprising of three species or more was considered to study the complex ecological community [9–13]. Furthermore, due to the existence of different important factors in the environment, such as stage structure, refuge, cannibalism, harvesting, toxicant, and scavenger, different prey-predator models incorporating such factors have been proposed and studied in literature, see, for example, [14–17].

On the other hand, the spread of diseases between the individuals of species is an important subject for study, which is known as an epidemiology, due to the existence of a variety of diseases in the environment. It is well known that after the SIR epidemic model of Kermack–Mckendrick [18], the field of epidemiology has received a lot of attention from the researchers, see [19] and the references therein. Recently, Zhou and Yao [20] proposed and studied a host-vector epidemic model with a stage structure for the vector. Further articles have been done also in the field of epidemiology in one field that is called ecoepidemiology, which has become a major field of study, see, for example, [21–24] and the references therein. Recently, Wang et al. [25] proposed and studied

a generalized ecoepidemiological system with prey refuge. The saturation incidence kinetics and a generalized functional response are used to describe the contact process and the predation process, respectively. They show that the prey refuge can control the spread of disease by the relative level of prey refuge.

Keeping the above studies in view, an ecoepidemiological model is proposed consisting of prey-predator incorporating a prey refuge, infectious disease of SI type in prey species only, and the predator preys upon the prey (susceptible and infected) according to the modified Holling type-II functional responses. It is assumed further that there is a harvesting process on the predator species due to external resources.

2. Model Formulation

In this section, a prey-predator system incorporating disease in a prey is formulated mathematically. It is assumed that there is a refuge for prey to protect themselves against predation and the harvesting process on predator. The following hypotheses are adopted in constricting the mathematical model that describes the dynamics of the above real-world ecoepidemiological system.

- (1) In the presence of infectious disease, the prey population is divided into two compartments, namely, susceptible prey that is denoted to their density at time *t* by S(T) and infected prey, which is denoted to their density at time *T* by I(T). Therefore, at time *T*, the total prey population density is X(T) = S(T) + I(T)
- (2) It is assumed that in the absence of predator that is denoted to their density at time *T* by *Y*(*T*), the susceptible prey grows logistically with carrying a capacity *K* > 0 and an intrinsic growth rate *r* > 0, while the infected prey cannot reproduce and rather than that, it competes for food and space
- (3) The availability of refuge for the prey in the environment keeps m ∈ (0, 1) portion of them safe from predation and left (1 m) from the total density of the prey facing predator
- (4) The disease transfers from infected to susceptible individuals by contact with an infected rate given by β > 0
- (5) There is a harvesting process on the predator, due to an external force, represented by a rate $(l_1E)/(l_2E + l_3Y)$, where E > 0 is the effort applied to harvest the prey; $l_1 > 0$ is the catchability coefficient while $l_2 > 0$ and $l_3 > 0$ are suitable constants
- (6) The predator consumed the prey according to the modified Holling type-II with attack rates $a_1 > 0$ and $a_2 > 0$ for susceptible prey and infected prey, respectively, half saturation constant b > 0, and infected prey's preference rate c > 0, while the conversion rate is given by e > 0

(7) Finally, there is a disease death rate given by $d_1 > 0$ and the natural death rate for predator $d_2 > 0$

According to the above hypotheses, the dynamics of the above described ecoepidemiological real-world system can be represented by the following set of differential equations.

$$\begin{aligned} \frac{dS}{dT} &= rS\left(1 - \frac{S+I}{K}\right) - \beta SI - \frac{a_1(1-m)SI}{b+(1-m)S+c(1-m)I},\\ \frac{dI}{dT} &= \beta SI - \frac{a_2(1-m)IY}{b+(1-m)S+c(1-m)I} - d_1I,\\ \frac{dY}{dT} &= \frac{e(a_1S-a_2I)(1-m)Y}{b+(1-m)S+c(1-m)I} + d_2Y - \frac{l_1EY}{l_2E+l_3Y}, \end{aligned}$$
(1)

with initial data $S(0) \ge 0$, $I(0) \ge 0$, and $Y(0) \ge 0$. Clearly, the interaction functions in system (1) are continuously differential functions on the positive octant.

Hence, they are Lipschitzian. Therefore, the solution of system (1) exits and is unique. Now, for further simplifying the above system, the following dimensionless variables and constants are used.

$$rT = t,$$

$$\frac{S}{K} = x,$$

$$\frac{I}{K} = y,$$

$$\frac{a_1Y}{rK} = z,$$

$$w_1 = \frac{\beta K}{r},$$

$$w_2 = \frac{\beta K}{r},$$

$$w_3 = \frac{a_2}{a_1}$$

$$w_4 = \frac{d_1}{r},$$

$$w_5 = \frac{ea_1}{r},$$

$$w_6 = \frac{d}{r},$$

$$w_7 = \frac{l_1a_1E}{l_3r^2K},$$

$$w_8 = \frac{l_2Ea_1}{l_3rK}$$

(2)

Thus, system (1) can be turned into the following dimensionless form:

$$\begin{split} \frac{dx}{dt} &= x \left[1 - x - y - w_1 y - \frac{(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} \right] \\ &= x f_1(x, y, z), \\ \frac{dy}{dt} &= y \left[w_1 x - \frac{w_3(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} - w_4 \right] \\ &= y f_2(x, y, z), \end{split}$$

$$\begin{aligned} \frac{dz}{dt} &= z \left[\frac{w_5(1-m)(x+w_3y)}{w_2 + (1-m)x + c(1-m)y} - w_6 - \frac{w_7}{w_8 + z} \right] \\ &= z f_3(x,y,z), \end{aligned}$$
(3)

with initial data (x(0), y(0), z(0)) belonging to $\mathbb{R}^4_+ = \{(x, y, z) \in \mathbb{R}^3 : x(t) > 0, y(t) > 0, z(t) > 0\}$. Clearly, according to the biological meaning of the parameters, we have always that $w_5 < 1$. Moreover, all solutions of system (3) initiated in \mathbb{R}^4_+ are uniformly bounded as shown in the following theorem.

Theorem 1. All the solutions of system (3), which initiate in \mathbb{R}^3_+ , are uniformly bounded.

Proof. According to the first equation of system (3) with the help of the Gronwall lemma, we have x(t) < 1. Let (x(t), y(t), z(t)) be any solution of system (3) initiated in \mathbb{R}^3_+ , and let W be a function defined by W(t) = x(t) + y(t) + z(t). Then, by differentiating this function with respect to time, it gives

$$\frac{dW}{dt} = \frac{dx}{dt} + \frac{dy}{dt} + \frac{dz}{dt} \le 2 - \mu_1 W, \tag{4}$$

where $\mu_1 = \min \{1, w_4, w_6\}$. Therefore, it obtained

$$\frac{dW}{dt} + \mu_1 W \le 2. \tag{5}$$

Again, using the Gronwall lemma gives that as $t \longrightarrow \infty$, $W(t) \le (2)/(\mu_1)$. Thus, all the solutions of system (1) are uniformly bounded.

3. Stability and Persistence

In this section, the stability analysis of all possible equilibrium points of system (3) is investigated. The persistence conditions of system (3) are established. Now, straightforward computation shows that the system (1) has at most five nonnegative equilibrium points. The existence conditions for each of them can be summarized as follows: The trivial equilibrium point (TEP), say $E_0 = (0, 0, 0)$ always exists. The disease-predator-free equilibrium point (DPFEP), namely, $E_1 = (x', 0, 0) = (1, 0, 0)$ always exists. The predator-free equilibrium point (PFEP), namely, $E_2 = (\bar{x}, \bar{y}, 0) = (((w_4)/(w_1)), ((w_1 - w_4)/(w_1(1 + w_1))), 0)$ exists in the *xy*-plane provided that

$$w_4 < w_1. \tag{6}$$

The disease-free equilibrium point (DFEP) is denoted by $E_3 = (\hat{x}, 0, \hat{z})$ where

$$\widehat{z} = \frac{(1-\widehat{x})(w_2 + (1-m)\widehat{x})}{(1-m)},\tag{7}$$

while \hat{x} is a positive root of the following third-order equation:

$$\alpha_1 \hat{x}^3 + \alpha_2 \hat{x}^2 + \alpha_3 \hat{x} + \alpha_4 = 0, \qquad (8)$$

with

$$\begin{aligned} &\alpha_1 = N^2(w_6 - w_5), \\ &\alpha_2 = N^2(w_5 - w_6) + w_2 N(2w_6 - w_5), \\ &\alpha_3 = w_8 N^2(w_5 - w_6) + w_2 N(w_5 - 2w_6) + w_2^2 w_6 - N^2 w_7, \\ &\alpha_4 = -w_2 N(w_6 w_8 + w_7) - w_2^2 w_6 < 0. \end{aligned}$$
(9)

Hence, E_3 exists uniquely in the positive quadrant of xz-plane if one of the following sets of conditions holds.

$$\alpha_1 > 0 \text{ and } \alpha_2 > 0$$

$$\alpha_1 > 0 \text{ and } \alpha_3 < 0$$
(10)

The coexistence equilibrium point (CEP) is denoted by $E_4 = (x^*, y^*, z^*)$, where

$$y^{*} = \frac{w_{3} + w_{4} - (w_{1} + w_{3})x^{*}}{w_{3}(1 + w_{1})},$$

$$z^{*} = \frac{(w_{1}x^{*} - w_{4})(w_{2} + Nx^{*} + cNy^{*})}{Nw_{3}},$$
(11)

while x^* is a positive root of the following third-order equation.

$$\sigma_1(x^*)^3 + \sigma_2(x^*)^2 + \sigma_3 x^* + \sigma_4 = 0.$$
 (12)

Here,

$$\sigma_{1} = \gamma_{1}(\theta_{1} - \theta_{3}w_{6}),$$

$$\sigma_{2} = \theta_{1}\gamma_{2} + \theta_{2}\gamma_{1} - \theta_{3}\gamma_{2}w_{6} - \theta_{4}\gamma_{1}w_{6},$$

$$\sigma_{3} = \theta_{1}\theta_{6} + \theta_{2}\gamma_{2} - \theta_{3}\theta_{6}w_{6} - \theta_{4}\gamma_{2}w_{6} - \theta_{5}\theta_{3},$$

$$\sigma_{4} = \theta_{2}\theta_{6} - \theta_{4}\theta_{6}w_{6} - \theta_{3}\theta_{4},$$
(13)

with

$$\begin{split} \theta_{1} &= w_{3}w_{5}N(1-w_{3}), \\ \theta_{2} &= w_{3}w_{5}N(w_{3}+w_{4}), \\ \theta_{3} &= N(w_{3}(1+w_{1})-c(w_{1}+w_{3})), \\ \theta_{4} &= w_{2}w_{3}(1+w_{1})+cN(w_{3}+w_{4}), \\ \theta_{5} &= w_{7}w_{3}^{2}N(1+w_{1}), \\ \theta_{6} &= \theta_{3}+w_{8}Nw_{3}^{2}(1+w_{1}), \\ \gamma_{1} &= w_{1}N[w_{3}(1+w_{3})-c(w_{1}+w_{3})], \\ \gamma_{2} &= w_{3}(1+w_{1})(w_{1}w_{2}-w_{4}N) \\ &+ cNw_{1}(w_{3}+w_{4})+cNw_{4}(w_{1}+w_{3}). \end{split}$$
(14)

Clearly, the CEP that is denoted by E_4 exists uniquely in the interior of the positive octant provided that the following condition holds.

$$\frac{w_4}{w_1} < x^* < \frac{w_3 + w_4}{w_1 + w_3},\tag{15}$$

with one set of the following sets of conditions

$$\left.\begin{array}{l} \sigma_{1} > 0, \sigma_{4} < 0, \sigma_{2} > 0\\ \sigma_{1} > 0, \sigma_{4} < 0, \sigma_{3} < 0\\ \sigma_{1} < 0, \sigma_{4} > 0, \sigma_{2} < 0\\ \sigma_{1} < 0, \sigma_{4} > 0, \sigma_{3} > 0 \end{array}\right\}$$
(16)

In the following, the local stability analysis of the above equilibrium points is carried out. The Jacobian matrix of system (1) at point (x, y, z) can be written as

$$J(x, y, z) = \left(u_{ij}\right)_{3\times 3} \tag{17}$$

Here,

$$\begin{split} u_{11} &= -x + \frac{(1-m)^2 xz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2} + f_1(x, y, z), \\ u_{12} &= -x - w_1 x + \frac{c(1-m)^2 xz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2}, \\ u_{13} &= \frac{-(1-m)x}{\left[w_2 + (1-m)x + c(1-m)y\right]}, \\ u_{21} &= w_1 y + \frac{w_3(1-m)^2 yz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2}, \end{split}$$

$$u_{22} = \frac{w_3(1-m)^2 yz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2} + f_2(x, y, z),$$

$$u_{23} = \frac{-w_3(1-m)y}{\left[w_2 + (1-m)x + c(1-m)y\right]},$$

$$u_{31} = \frac{w_2 w_5(1-m)z + w_5(c-w_3)(1-m)^2 yz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2},$$

$$u_{32} = \frac{w_2 w_3 w_5(1-m)z + w_5(w_3-c)(1-m)^2 xz}{\left[w_2 + (1-m)x + c(1-m)y\right]^2},$$

$$u_{33} = \frac{w_7 z}{\left[w_8 + z\right]^2} + f_3(x, y, z).$$
(18)

Therefore, the eigenvalues of the Jacobian matrix at E_0 are

$$\lambda_{0x} = 1 > 0,$$

$$\lambda_{0y} = -w_4 < 0,$$

$$\lambda_{0z} = -w_6 - \frac{w_7}{w_8} < 0.$$
(19)

Hence, the TEP is a saddle point with unstable manifold in *x*-direction and stable manifold in the *yz*-plane.

The Jacobian matrix of system (3) at the DPFEP that is represented by $E_1 = (1, 0, 0)$ is given by

$$J(E_1) = \begin{pmatrix} -1 & -1 - w_1 & -\frac{(1-m)}{w_2 + (1-m)} \\ 0 & w_1 - w_4 & 0 \\ 0 & 0 & \frac{w_5(1-m)}{w_2 + (1-m)} - w_6 - \frac{w_7}{w_8} \end{pmatrix}$$
(20)

Therefore, the eigenvalues are written as

$$\lambda_{1x} = -1 < 0,$$

$$\lambda_{1y} = w_1 - w_4,$$

$$\lambda_{1z} = \frac{w_5(1-m)}{w_2 + (1-m)} - w_6 - \frac{w_7}{w_8}.$$
(21)

It is clear that all the eigenvalues are negative and then, the DPFEP is locally asymptotically stable provided that the following conditions are hold.

$$w_1 < w_4, \tag{22a}$$

$$\frac{w_5(1-m)}{w_2+(1-m)} < w_6 + \frac{w_7}{w_8},$$
(22b)

The Jacobian matrix of system (3) at the PFEP is denoted by $E_2 = (\bar{x}, \bar{y}, 0)$.

$$J(E_2) = \begin{pmatrix} -\bar{x} & -(1+w_1)\bar{x} & -\frac{(1-m)\bar{x}}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} \\ w_1\bar{y} & 0 & -\frac{w_3(1-m)\bar{y}}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} \\ 0 & 0 & \frac{w_5(1-m)(\bar{x}+w_3\bar{y})}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} - w_6 - \frac{w_7}{w_8} \end{pmatrix}$$
(23)

Clearly, the characteristic equation of $J(E_2)$ can be written as

$$\begin{bmatrix} \lambda_2^2 + \bar{x}\lambda_2 + w_1(1+w_1)\bar{x}\bar{y} \end{bmatrix} \\ \cdot \begin{bmatrix} w_5(1-m)(\bar{x}+w_3\bar{y}) \\ w_2 + (1-m)\bar{x} + c(1-m)\bar{y} \end{bmatrix} - w_6 - \frac{w_7}{w_8} - \lambda_2 \end{bmatrix} = 0.$$
(24)

Accordingly, all the eigenvalues of $J(E_2)$ have negative real parts and then, the PFEP is locally asymptotically stable provided that the following condition holds

$$\frac{w_5(1-m)(\bar{x}+w_3\bar{y})}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} < w_6 + \frac{w_7}{w_8}.$$
 (25)

The Jacobian matrix of system (3) at the DFEP that is given by E_3 can be written as

$$J(E_3) = \left(b_{ij}\right)_{3\times 3},\tag{26}$$

where

$$b_{11} = \left[-1 + \frac{(1-m)^2 \hat{z}}{(w_2 + (1-m)\hat{x})^2} \right] \hat{x},$$

$$b_{12} = \left[-1 - w_1 + \frac{(1-m)^2 \hat{z}}{(w_2 + (1-m)\hat{x})^2} \right] \hat{x},$$

$$b_{13} = -\frac{(1-m)\hat{x}}{w_2 + (1-m)\hat{x}},$$

$$b_{21} = b_{23} = 0,$$

$$b_{22} = w_1 \hat{x} - \frac{w_3(1-m)\hat{z}}{w_2 + (1-m)\hat{x}} - w_4,$$

$$b_{31} = \frac{w_2 w_5 (1-m)\hat{z}}{(w_2 + (1-m)\hat{x})^2},$$

$$b_{32} = \left[\frac{w_2 w_3 w_5 (1-m) + w_5 (1-m)^2 (w_3 - c) \hat{x}}{(w_2 + (1-m)\hat{x})^2} \right] \hat{z},$$

$$b_{33} = \frac{w_7 \hat{z}}{(w_8 + \hat{z})^2}.$$
(27)

Hence, the characteristic equation of $J(E_3)$ can be written as

$$\left[\lambda_{3}^{2} - (b_{11} + b_{33})\lambda_{3} + (b_{11}b_{33} - b_{13}b_{31})\right][b_{22} - \lambda_{3}] = 0.$$
(28)

Straightforward computations show that all the roots (eigenvalues) of the characteristic equation (28) have negative real parts, and then, the DFEP is locally asymptotically stable provided that the following conditions hold.

$$\frac{(1-m)^2 \hat{x} \hat{z}}{(w_2 + (1-m)\hat{x})^2} + \frac{w_7 \hat{z}}{(w_8 + \hat{z})^2} < \hat{x},$$
 (29a)

$$\frac{w_7}{(w_8 + \hat{z})^2} < \frac{w_7 (1 - m)^2 \hat{z}}{(w_8 + \hat{z})^2 (w_2 + (1 - m) \hat{x})^2} + \frac{w_2 w_5 (1 - m)^2}{(w_2 + (1 - m) \hat{x})^3},$$
(29b)

$$w_1 \hat{x} < \frac{w_3 (1-m)\hat{z}}{w_2 + (1-m)\hat{x}} + w_4.$$
(29c)

Now, the local stability conditions for the CEP are established in the following theorem.

Theorem 2. Suppose that the CEP exists; then, it is locally asymptotically stable in the \mathbb{R}^3_+ if the following sufficient conditions hold.

$$\frac{(1-m)^2 x^* z^* + w_3 (1-m)^2 y^* z^*}{\left[w_2 + (1-m)x^* + c(1-m)y^*\right]^2} + \frac{w_7 z^*}{\left[w_8 + z^*\right]^2} < x^*, \quad (30a)$$

$$\frac{c(1-m)^2 z^*}{[w_2 + (1-m)x^* + c(1-m)y^*]^2} < 1 + w_1,$$
(30b)

$$\begin{array}{c} w_{3} \leq c < \frac{w_{2}w_{3} + w_{3}(1-m)x^{*}}{(1-m)x^{*}} \\ or \\ w_{c}(1-m)y^{*} - w_{c} \end{array} \right\},$$
(30c)

$$\frac{\omega_3(1-m)y-\omega_2}{(1-m)y^*} < c \le w_3$$

$$\frac{a_{13}a_{31}}{a_{33}} < a_{11} < \min\left\{\frac{a_{12}a_{21}}{a_{22}}, \frac{a_{12}a_{31}}{a_{32}}\right\},\tag{30d}$$

$$a_{22} < \frac{a_{21}a_{32}}{a_{31}},$$
 (30e)

$$M_1 > M_2$$
, (30f)

where all symbols are clearly described in the proof.

Proof. According to Equation (17), the Jacobian matrix at E_4 is given by

$$J(E_4) = \left(a_{ij}\right)_{3\times 3},\tag{31}$$

with $a_{ij} = u_{ij}(x^*, y^*, z^*)$ for all i, j = 1, 2, 3. Therefore, the characteristic equation of $J(E_4)$ can be written as

$$\lambda_4^{\ 3} + D_1 \lambda_4^{\ 2} + D_2 \lambda_4 + D_3 = 0, \tag{32}$$

where

$$D_{1} = -[a_{11} + a_{22} + a_{33}],$$

$$D_{2} = a_{11}a_{22} - a_{12}a_{21} + a_{11}a_{33} - a_{13}a_{31} + a_{22}a_{33} - a_{23}a_{32},$$

$$D_{3} = -[a_{33}(a_{11}a_{22} - a_{12}a_{21}) + a_{23}(a_{12}a_{31} - a_{11}a_{32}) + a_{13}(a_{21}a_{32} - a_{22}a_{31})].$$
(33)

Moreover, it is easy to verify that

$$D_1D_2 - D_3 = -(a_{11} + a_{22})(a_{11}a_{22} - a_{12}a_{21}) - (a_{11} + a_{33})(a_{11}a_{33} - a_{13}a_{31}) - (a_{22} + a_{33})(a_{22}a_{33} - a_{23}a_{32}) - 2a_{11}a_{22}a_{33} + a_{12}a_{23}a_{31} + a_{13}a_{21}a_{32} = M_1 - M_2,$$
(34)

where

$$M_{1} = -(a_{11} + a_{33})(a_{11}a_{33} - a_{13}a_{31}) - 2a_{11}a_{22}a_{33} + a_{12}a_{23}a_{31} > 0, M_{2} = (a_{11} + a_{22})(a_{11}a_{22} - a_{12}a_{21})$$
(35)

$$+ (a_{22} + a_{33})(a_{22}a_{33} - a_{23}a_{32}) - a_{13}a_{21}a_{32} > 0.$$

Recall that according to Routh-Hurwitz criterion, Equation (32) has roots (eigenvalues) with negative real parts if and only if $D_1 > 0$, $D_3 > 0$ and $D_1D_2 - D_3 > 0$. Now, straightforward computation shows that all the requirements of the Routh-Hurwitz criterion are satisfied provided that the conditions (30a), (30b), (30c), (30e), and (30f) hold. Hence, the CEP is locally asymptotically stable and the proof is complete.

It is well known that the persistence of an ecological system is an important subject for both environment and ecologists. The coexistence of all the species along the entire time life is known as persistence. Mathematically, this is equivalent to unapproachable of the trajectory to any of the boundary planes of axis. It is well known that system (3) has two subsystems. The first subsystem exists in the *x y*-plane, in case in the absence of the predator, and can be represented in system (36). However, the second subsystem exists in the xz-plane, in the absence of disease, and can be represented in system (37).

$$\frac{dx}{dt} = x[1 - x - y - w_1 y] = g_1(x, y),$$

$$\frac{dy}{dt} = y[w_1 x - w_4] = g_2(x, y),$$
(36)

$$\begin{aligned} \frac{dx}{dt} &= x \left[1 - x - \frac{(1 - m)z}{w_2 + (1 - m)x} \right] = g_3(x, z), \\ \frac{dz}{dt} &= z \left[\frac{w_5(1 - m)x}{w_2 + (1 - m)x} - w_6 - \frac{w_7}{w_8 + z} \right] = g_4(x, z). \end{aligned}$$
(37)

Recall that according to Bendixson-Dulac Criteria [26], the planar system X' = G(X), where $X = (x_1, x_2)^T$ and $G = (g, f)^T$ has no periodic dynamics provided that there is a C^1 function $\varphi(x_1, x_2)$ (known as Dulac function) so that the expression $(\partial(\varphi g))/(\partial x) + (\partial(\varphi f))/(\partial y)$ has the same sign and does not equal zero almost everywhere in a simply connected region of the plane.

Consequently, by choosing the Dulac function (1)/(xy), it is easy to verify that system (36) has no periodic dynamics in the interior of *xy*-plane, while with the Dulac function (1)/(xz), system (37) has no periodic dynamics fall in the interior of *xz*-plane provided that the following condition holds.

$$\frac{(1-m)^2}{\left[w_2 + (1-m)x\right]^2} + \frac{w_7}{x\left[w_8 + z\right]^2} \neq \frac{1}{z}.$$
 (38)

Now, in order to established the persistence conditions for system (3), the average Lyapunov function will be used as shown in the following theorem.

Theorem 3. Assuming that there is no periodic dynamics in the boundary planes; then, system (3) is persistent provided that the following conditions hold.

$$w_1 > w_4, \tag{39a}$$

$$\frac{w_5(1-m)}{w_2+(1-m)} > w_6 + \frac{w_7}{w_8},$$
(39b)

$$\frac{w_5(1-m)(\bar{x}+w_3\bar{y})}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} > w_6 + \frac{w_7}{w_8},$$
(39c)

$$w_1 \hat{x} > \frac{w_3 (1-m)\hat{z}}{w_2 + (1-m)\hat{x}} + w_4.$$
 (39d)

Proof. Consider the following function $P(x, y, z) = x^{q_1}y^{q_2}z^{q_3}$, where q_i , $\forall i = 1, 2, 3$ are positive constants to be determined. Then, we obtain

$$\begin{split} \Lambda &= \frac{P'}{P} = q_1 \left[1 - x - y - w_1 y - \frac{(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} \right] \\ &+ q_2 \left[w_1 x - \frac{w_3(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} - w_4 \right] \\ &+ q_3 \left[\frac{w_5(1 - m)(x + w_3 y)}{w_2 + (1 - m)x + c(1 - m)y} - w_6 - \frac{w_7}{w_8 + z} \right]. \end{split}$$

$$\end{split} \tag{40}$$

Now, according to the average Lyapunov function, the proof is done if we could prove that this function Λ is positive at each boundary equilibrium point. Note that at E_0 , it obtains

$$\Lambda(E_0) = q_1 - w_4 q_2 - \left(w_6 + \frac{w_7}{w_8}\right) q_3.$$
(41)

So, for any positive constants q_i , $\forall i = 1, 2, 3$ with q_1 sufficiently large compared with q_2 and q_3 , we get $\Lambda(E_0) > 0$. At the DPFEP that is represented by $E_1 = (1, 0, 0)$, we have

$$\Lambda(E_1) = (w_1 - w_4)q_2 + \left[\frac{w_5(1 - m)}{w_2 + (1 - m)} - \left(w_6 + \frac{w_7}{w_8}\right)\right]q_3.$$
(42)

Clearly, $\Lambda(E_1) > 0$ under the conditions (39a) and (39b). Now, at the PFEP that is denoted by $E_2 = (\bar{x}, \bar{y}, 0)$, it obtains

$$\Lambda(E_2) = q_3 \left[\frac{w_5(1-m)(\bar{x}+w_3\bar{y})}{w_2 + (1-m)\bar{x} + c(1-m)\bar{y}} - \left(w_6 + \frac{w_7}{w_8}\right) \right].$$
(43)

Here, $\Lambda(E_2) > 0$ under condition (39c). Finally, at the DFEP that is given by $E_3 = (\hat{x}, 0, \hat{z})$, we have

$$\Lambda(E_3) = q_2 \left[w_1 \hat{x} - \frac{w_3 (1-m) \hat{z}}{w_2 + (1-m) \hat{x}} - w_4 \right].$$
(44)

Clearly, condition (39d) guarantees that $\Lambda(E_3) > 0$. Hence, the system has no omega limit sets on the boundary planes and thus, the system is persists.

Now, the global dynamics of system (3) is investigated using suitable Lyapunov functions as shown in the following theorem.

Theorem 4. Assuming that the DPFEP which is given by $E_1 = (1, 0, 0)$ is locally asymptotically stable, then, it is a globally asymptotically stable provided that

$$1 + w_1 < w_4, \tag{45a}$$

$$\frac{1-m}{w_2} < w_6. \tag{45b}$$

Proof. Consider the positive definite function

$$L_1 = x - 1 - \ln(x) + y + z.$$
(46)

It is clear that $L_1 : \mathbb{R}^3_+ \longrightarrow \mathbb{R}$ and $L_1(1, 0, 0) = 0$ with $L_1(x, y, z) > 0$, for all $(x, y, z) \neq (1, 0, 0)$. Now the derivative of L_1 with respect to time can be written, after some algebraic computation, as

$$\frac{dL_1}{dt} < -(x-1)^2 - [w_4 - (1+w_1)]y - \left[w_6 - \frac{(1-m)}{w_2}\right]z.$$
(47)

Therefore, by using conditions (45a) and (45b), it obtains $(dL_1)/(dt)$ which is negatively defined in \mathbb{R}^3_+ . Thus, according to Lyapunov's second theorem, E_1 is globally asymptotically stable.

Theorem 5. Assuming that the PFEP that is given by is locally asymptotically stable, then, it is globally asymptotically stable in \mathbb{R}^3_+ provided that

$$\left(\frac{1+w_1}{w_1}\right)\frac{w_3(1-m)\bar{y}}{w_2} + \frac{(1-m)\bar{x}}{w_2} < \frac{w_6}{w_5}.$$
 (48)

Proof. Consider the positive definite function

$$L_{2} = C_{1} \left[x - \bar{x} - \bar{x} \ln \left(\frac{x}{\bar{x}} \right) \right] + C_{2} \left[y - \bar{y} - \bar{y} \ln \left(\frac{y}{\bar{y}} \right) \right] + C_{2} z.$$

$$(49)$$

It is clear that $L_2 : \mathbb{R}^3_+ \longrightarrow \mathbb{R}$ and $L_2(\bar{x}, \bar{y}, 0) = 0$ with $L_2(x, y, z) > 0$, for all $(x, y, z) \neq (\bar{x}, \bar{y}, 0) = (((w_4)/(w_1)), ((w_1 - w_4)/(w_1(1 + w_1))), 0)$. Then, the derivative of L_2 with respect to time can be written, after some algebraic steps as

$$\begin{aligned} \frac{dL_2}{dt} &= -C_1 (x - \bar{x})^2 - [C_1 (1 + w_1) - C_2 w_1] (x - \bar{x}) (y - \bar{y}) \\ &- [C_1 - C_2 w_5] \frac{(1 - m)xz}{w_2 + (1 - m)x + c(1 - m)y} \\ &- [C_2 - C_3 w_5] \frac{w_3 (1 - m)yz}{w_2 + (1 - m)x + c(1 - m)y} \\ &- \left[C_3 w_6 - C_2 \frac{w_3 (1 - m)\bar{y}}{w_2} - C_1 \frac{(1 - m)\bar{x}}{w_2} \right] z \\ &- C_3 \frac{w_7 z}{w_8 + z}. \end{aligned}$$
(50)

So by choosing the positive constants as $C_1 = 1$, $C_2 = ((1 + w_1)/(w_1))$, and $C_3 = (1/w_5)$, the following is obtained

$$\frac{dL_2}{dt} = -(x-\bar{x})^2 - \left(\frac{1}{w_1}\right) \frac{w_3(1-m)yz}{w_2 + (1-m)x + c(1-m)y} \\
- \left[\frac{w_6}{w_5} - \left(\frac{1+w_1}{w_1}\right) \frac{w_3(1-m)\bar{y}}{w_2} - \frac{(1-m)\bar{x}}{w_2}\right]z \\
- \left(\frac{1}{w_5}\right) \frac{w_7z}{w_8 + z}, \\
\frac{dL_2}{dt} = -(x-\bar{x})^2 - \left(\frac{1}{w_1}\right) \frac{w_3(1-m)yz}{w_2 + (1-m)x + c(1-m)y}.$$
(51)

Clearly, under the condition (48), (dL_2/dt) is negatively semidefinite, because $(dL_2/dt) = 0$ if and only if $x = \bar{x}$ and z = 0 with any value of y. Therefore, according to Lyapunov's second theorem, E_2 is a stable point.

Now, form the first equation of system (3), we have $(dx/dt) = \bar{x}[1 - \bar{x} - (1 + w_1)y]$, which is equal to zero if and only if $y = \bar{y}$. Hence, the set $\{E_2\}$ is the largest invariant set for which $(dL_2/dt) = 0$. Accordingly, the global asymptotic stability of the PFEP follows from LaSalle's invariance principle. This completes the proof.

Theorem 6. Assuming that the DFEP that is given by $E_3 = (\hat{x}, 0, \hat{z})$ is locally asymptotically stable, then, it is globally asymptotically stable in \mathbb{R}^3_+ provided that the following sufficient conditions hold

$$\frac{(1-m)^2 \hat{z}}{w_2 \hat{N}_1} < 1,$$
(52a)

$$c(1-m)^2 w_4 \hat{z} + (1+w_1) w_2^2 \hat{N}_1 \hat{x} < (1+w_1) w_2 w_4 \hat{N}_1,$$
(52b)

$$p_{13}^{2} < 4p_{11}p_{33}, \tag{52c}$$

$$2p_{33}(z-\hat{z})^2 < \left[\sqrt{p_{11}}(x-\hat{x}) + \sqrt{p_{33}}(z-\hat{z})\right]^2,$$
(52d)

where all symbols are given in the proof.

Proof. Consider the positive definite function

$$L_{3} = \rho_{1} \left[x - \hat{x} - \hat{x} \ln \left(\frac{x}{\hat{x}} \right) \right] + \rho_{2} y + \rho_{3} \left[z - \hat{z} - \hat{z} \ln \left(\frac{z}{\hat{z}} \right) \right].$$
(53)

It is clear that $L_3 : \mathbb{R}^3_+ \longrightarrow \mathbb{R}$ and $L_3(\hat{x}, 0, \hat{z}) = 0$ with $L_3(x, y, z) > 0$, for all $(x, y, z) \neq (\hat{x}, 0, \hat{z})$. Then, the derivative of L_3 with respect to time can be written, after some algebraic steps as

$$\begin{aligned} \frac{dL_{3}}{dt} &\leq -\rho_{1} \left[1 - \frac{(1-m)^{2} \hat{z}}{w_{2} \hat{N}_{1}} \right] (x-\hat{x})^{2} - \frac{1-m}{N_{1} \hat{N}_{1}} \\ & \cdot \left[\rho_{1} \hat{N}_{1} - \rho_{3} w_{2} w_{5} \right] (x-\hat{x}) (z-\hat{z}) + \rho_{3} \frac{w_{7}}{w_{8} \hat{N}_{2}} (z-\hat{z})^{2} \\ & - \left[\rho_{1} \left(\frac{(1+w_{1}) w_{2} \hat{N}_{1} - c(1-m)^{2} \hat{z}}{w_{2} \hat{N}_{1}} \right) - \rho_{2} w_{2} \right] xy \\ & - \left[\rho_{2} w_{4} - \rho_{1} (1+w_{1}) \hat{x} \right] y - \frac{(1-m)^{2}}{N_{1} \hat{N}_{1}} \left[\rho_{1} c - \rho_{3} w_{5} \right] \hat{x} y \hat{z} \\ & - \frac{1-m}{N_{1}} \left[\rho_{2} w_{3} + \rho_{3} \frac{w_{5} (1-m) \hat{x}}{\hat{N}_{1}} \right] yz. \end{aligned}$$

$$(54)$$

where $N_1 = w_2 + (1 - m)x + c(1 - m)y$, $\hat{N}_1 = w_2 + (1 - m)\hat{x}$, $N_2 = w_8 + z$, and $\hat{N}_2 = w_8 + \hat{z}$. So, by choosing the positive constants as $\rho_1 = 1$, $\rho_2 = (((1 + w_1)w_2\hat{N}_1 - c(1 - m)^2\hat{z})/(w_2^2\hat{N}_1))$, which is positive under (52a), and $\rho_3 = (c/w_5)$. Therefore, we obtain

$$\frac{dL_{3}}{dt} < -\left[1 - \frac{(1-m)^{2}\hat{z}}{w_{2}\hat{N}_{1}}\right](x-\hat{x})^{2} - \frac{1-m}{N_{1}\hat{N}_{1}}\left[\hat{N}_{1} - cw_{2}\right](x-\hat{x}) \\
\cdot (z-\hat{z}) - \frac{cw_{7}}{w_{5}w_{8}\hat{N}_{2}}(z-\hat{z})^{2} + 2\frac{cw_{7}}{w_{5}w_{8}\hat{N}_{2}}(z-\hat{z})^{2} \\
- \left[\frac{(1+w_{1})w_{2}\hat{N}_{1} - c(1-m)^{2}\hat{z}}{w_{2}^{2}\hat{N}_{1}}w_{4} - (1+w_{1})\hat{x}\right]y.$$
(55)

Therefore, according to the sufficient conditions (52a)-(52c), it is easy to obtain

$$\begin{aligned} \frac{dL_3}{dt} &< -\left[\sqrt{p_{11}}(x-\hat{x}) + \sqrt{p_{33}}(z-\hat{z})\right]^2 + 2p_{33}(z-\hat{z})^2 \\ &- \left[\frac{(1+w_1)w_2\hat{N}_1 - c(1-m)^2\hat{z}}{w_2^2\hat{N}_1}w_4 - (1+w_1)\hat{x}\right]y. \end{aligned}$$
(56)

Here, $p_{11} = 1 - (((1-m)^2 \hat{z})/(w_2 \hat{N}_1)), p_{13} = ((1-m)/(N_1 \hat{N}_1))[\hat{N}_1 - cw_2]$, and $p_{33} = ((cw_7)/(w_5 w_8 \hat{N}_2)).$

Theorem 7. Assuming that the positive equilibrium point $E_4 = (x^*, y^*, z^*)$ is locally asymptotically stable, then, all the points that satisfy the following conditions belong to the basin of attraction of E_4 .

$$\frac{(1-m)^2 z^*}{w_2 N_1^*} < 1,$$
(57a)

$$q_{12}^{2} < 2q_{11}q_{22},$$
 (57b)

$$q_{13}^{2} < 2q_{11}q_{33}, \tag{57c}$$

$$q_{23}^{2} < 16q_{22}q_{33}, \tag{57d}$$

$$2\left[\sqrt{q_{22}}(y-y^*) + \sqrt{q_{33}}(z-z^*)\right]^2 < \frac{1}{2}\left[\sqrt{q_{11}}(x-x^*) + \sqrt{2q_{22}}(y-y^*)\right]^2 + \frac{1}{2}\left[\sqrt{q_{11}}(x-x^*) + \sqrt{2q_{33}}(z-z^*)\right]^2.$$
(57e)

Proof. Consider the positive definite function

$$L_4 = \left[x - x^* - x^* \ln\left(\frac{x}{x^*}\right) \right] + \left[y - y^* - y^* \ln\left(\frac{y}{y^*}\right) \right]$$

+ $\left[z - z^* - z^* \ln\left(\frac{z}{z^*}\right) \right].$ (58)

It is clear that $L_4 : \mathbb{R}^3_+ \longrightarrow \mathbb{R}$ and $L_4(x^*, y^*, z^*) = 0$ with $L_4(x, y, z) > 0$, for all $(x, y, z) \neq (x^*, y^*, z^*)$ with x > 0, y > 0 and z > 0. Then, the derivative of L_4 with respect to time can be written, after some algebraic steps as

$$\begin{aligned} \frac{dL_4}{dt} &= (x - x^*) \left[-(x - x^*) - (1 + w_1)(y - y^*) \right. \\ &\quad - \frac{(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} \\ &\quad + \frac{(1 - m)z^*}{w_2 + (1 - m)x^* + c(1 - m)y^*} \right] \\ &\quad + (y - y^*) \left[w_1(x - x^*) \right. \\ &\quad - \frac{w_3(1 - m)z}{w_2 + (1 - m)x + c(1 - m)y} \\ &\quad + \frac{w_3(1 - m)z^*}{w_2 + (1 - m)x^* + c(1 - m)y^*} \right] \\ &\quad + (z - z^*) \left[-\frac{w_7}{w_8 + z} + \frac{w_7}{w_8 + z^*} \\ &\quad + \frac{w_5(1 - m)(x + w_3y)}{w_2 + (1 - m)x + c(1 - m)y} \\ &\quad - \frac{w_5(1 - m)(x^* + w_3y^*)}{w_2 + (1 - m)x^* + c(1 - m)y^*} \right]. \end{aligned}$$
(59)

Therefore, we obtain

$$\begin{aligned} \frac{dL_4}{dt} &\leq -\frac{q_{11}}{2} (x - x^*)^2 - q_{12} (x - x^*) (y - y^*) - q_{22} (y - y^*)^2 \\ &- \frac{q_{11}}{2} (x - x^*)^2 - q_{13} (x - x^*) (z - z^*) - q_{33} (z - z^*)^2 \\ &+ 2q_{22} (y - y^*)^2 - q_{23} (y - y^*) (z - z^*) + 2q_{33} (z - z^*)^2, \end{aligned}$$

$$(60)$$

where

$$\begin{split} q_{11} &= \left[1 - \frac{(1-m)^2 z^*}{w_2 N_1^*}\right], \\ q_{12} &= \left[1 - \frac{c(1-m)^2 z^*}{w_2 N_1^*} - \frac{w_3 (1-m)^2 z^*}{w_2 N_1^*}\right], \\ q_{13} &= \frac{(1-m)}{N_1} \left[1 - \frac{w_2 w_5}{N_1^*} - \frac{c w_5 (1-m) y^*}{N_1^*} + \frac{w_3 w_5 (1-m) y^*}{N_1^*}\right], \end{split} \tag{61}$$

$$\begin{aligned} q_{23} &= \frac{(1-m)}{N_1} \left[w_3 + \frac{c w_5 (1-m) x^*}{N_1^*} - \frac{w_2 w_3 w_5}{N_1^*} - \frac{w_3 w_5 (1-m) x^*}{N_1^*}\right], \\ &\quad - \frac{w_3 w_5 (1-m) x^*}{N_1^*}\right], \\ q_{22} &= \frac{c w_3 (1-m)^2 z^*}{w_2 N_1^*}, \\ q_{33} &= \frac{w_7}{w_8 N_2^*}, \end{split}$$

with $N_1 = w_2 + (1 - m)x + c(1 - m)y$, $N_1^* = w_2 + (1 - m)x^* + c(1 - m)y^*$, and $N_2^* = w_8 + z^*$.

Clearly according to the condition (57a), then, $q_{11} > 0$, while q_{22} and q_{33} are always positive. Therefore, by using the conditions (57b)–(57d), we obtain

$$\frac{dL_4}{dt} \leq -\frac{1}{2} \left[\sqrt{q_{11}} (x - x^*) + \sqrt{2q_{22}} (y - y^*) \right]^2
- \frac{1}{2} \left[\sqrt{q_{11}} (x - x^*) + \sqrt{2q_{33}} (z - z^*) \right]^2
+ 2 \left[\sqrt{q_{22}} (y - y^*) + \sqrt{q_{33}} (z - z^*) \right]^2.$$
(62)

The first tow terms are clearly negative while the third term is clearly positive; therefore, the sufficient condition (57e) guarantees the negative definite of the derivative. Hence, the solution starting with any point satisfying conditions (57a)–(57e) approaches asymptotically to E_4 . So these points are belonging to the basin of attraction.

4. Bifurcation Analysis

Now, in order to compute the second derivative of the Jacobian matrix, system (3) is rewritten in the vector form as follows:

 $dX/dt = F(X), \quad X = (x, y, z)^T,$ $F = (F_1, F_2, F_3)^T = (xf_1, yf_2, zf_3)^T.(50)$

Therefore, according to the Jacobian matrix of system (3) at the point (x, y, z) that is given by Equation (17), we obtain

$$D^{2}F(x, y, z)(V, V) = [m_{ij}]_{3 \times 1},$$
 (63)

where $V = (v_1, v_2, v_3)^T$ be any vector and

$$\begin{split} m_{11} &= -2\left(1 - \frac{(1-m)^2[w_2 + c(1-m)y]z}{N_1^3}\right)v_1^2 - 2\left(\frac{(1-m)[w_2 + c(1-m)y]}{N_1^2}\right)v_1v_3 \\ &\quad -2\left(1 + w_1 - \frac{cz(1-m)^2[w_2 - (1-m)x + c(1-m)y]}{N_1^3}\right)v_1v_2 - 2\left(\frac{c^2(1-m)^3xz}{N_1^3}\right)v_2^2 + 2\left(\frac{c(1-m)^2x}{N_1^2}\right)v_2v_3, \\ m_{21} &= -2\left(\frac{w_3(1-m)^3yz}{N_1^3}\right)v_1^2 + 2\left(\frac{w_3(1-m)^2y}{N_1^2}\right)v_1v_3 + 2\left(w_1 + \frac{w_3z(1-m)^2[w_2 + (1-m)x - c(1-m)y]}{N_1^3}\right)v_1v_2 \\ &\quad + 2\left(\frac{cw_3z(1-m)^2[w_2 + (1-m)x]}{N_1^3}\right)v_2^2 - 2\left(\frac{w_3(1-m)[w_2 + (1-m)x]}{N_1^2}\right)v_2v_3, \\ m_{31} &= -2\left(\frac{w_5(1-m)^2z[w_2 + (1-m)(c-w_3)y]}{N_1^3}\right)v_1^2 - 2\left(\frac{w_5(1-m)^2z[w_2(c+w_3) - (c-w_3)(1-m)x + c(c-w_3)(1-m)y]}{N_1^3}\right)v_1v_2 \\ &\quad - 2\left(\frac{cw_5(1-m)^2z[w_2w_3 - (1-m)(c-w_3)y]}{N_1^3}\right)v_2^2 + 2\left(\frac{w_2w_5(1-m) + w_5(1-m)^2(c-w_3)y}{N_1^2}\right)v_1v_3 \\ &\quad + 2\left(\frac{w_2w_3w_5(1-m) - w_5(1-m)^2(c-w_3)x}{N_1^2}\right)v_2v_3 + 2\frac{w_7w_8}{N_2^3}v_3^2. \end{split}$$

Consequently, the conditions guarantee the occurrence of local bifurcations around the nonhyperbolic equilibrium points that are established in the following theorems.

Theorem 8. Assume that condition (22b) holds. Letting the parameter w_1 pass through the value $w_1^* \equiv w_4$, then, system (3) undergoes a transcritical bifurcation but neither saddle node nor pitchfork bifurcation can occur

Proof. Note that when the parameter w_1 passes through the positive value $w_1^* \equiv w_4$, it is easy to verify that the equilibrium point E_1 becomes a nonhyperbolic point and the Jacobian matrix of system (3) at (E_1, w_1^*) can be written as

$$J_{1} = J(E_{1}, w_{1}^{*})$$

$$= \begin{pmatrix} -1 & -1 - w_{4} & -\frac{(1-m)}{w_{2} + (1-m)} \\ 0 & 0 & 0 \\ 0 & 0 & \frac{w_{5}(1-m)}{w_{2} + (1-m)} - w_{6} - \frac{w_{7}}{w_{8}} \end{pmatrix}.$$
(65)

Thus, J_1 has the following eigenvalues $\lambda_{11}^{[1]} = -1$, $\lambda_{12}^{[1]} = 0$, and $\lambda_{13}^{[1]} = (w_5(1-m)/w_2 + (1-m)) - w_6 - (w_7/w_8)$. Clearly, $\lambda_{13}^{[1]}$ is negative under the condition (22b).

Let $V_1 = (v_{11}, v_{12}, v_{13})^T$ be the eigenvector of J_1 corresponding to the eigenvalue $\lambda_{12}^{[1]} = 0$. Then, straightforward computation gives $V_1 = (\beta_1 v_{12}, v_{12}, 0)^T$, where $\beta_1 = -(1 + w_4)$; v_{12} represents any nonzero real number.

Also, let $\Psi_1 = (\psi_{11}, \psi_{12}, \psi_{13})^T$ represent the eigenvector of J_1^T that corresponds to the eigenvalue $\lambda_{12}^{[1]} = 0$. Then, direct calculation shows that $\Psi_1 = (0, \psi_{12}, 0)^T$, where ψ_{12} is any nonzero real number.

Since the partial derivative of vector field F with respect to the parameter w_1 is given by $\partial F/\partial w_1 = F_{w_1} = (-xy, xy, 0)^T$, hence, we obtain $F_{w_1}(E_1, w_1^*) = (0, 0, 0)^T$. Therefore,

$$\Psi_1^{T} [F_{w_1}(E_1, w_1^*)] = 0.$$
(66)

Thus system (3) at E_1 with $w_1 = w_1^*$ does not experience saddle-node bifurcation in view of the Sotomayor theorem [26]. Now, we have

$$\Psi_{1}^{T} \left[DF_{w_{1}}(E_{1}, w_{1}^{*}) V_{1} \right] = v_{12}\psi_{12} \neq 0,$$
(67)

where DF_{w_1} represents the derivative of F_{w_1} with respect to X. In addition, by substituting the value of E_1, w_1^*, V_1 , and Ψ_1 in Equation (63), then we get

$$\Psi_1^{T} \left[D^2 F(E_1, w_1^*)(V_1, V_1) \right] = 2\psi_{12} v_{12}^2 \beta_1 w_4 \neq 0.$$
 (68)

Accordingly, by the Sotomayor theorem, system (3) near the equilibrium point E_1 with $w_1 = w_1^*$ possesses a transcritical bifurcation, while pitchfork bifurcation cannot occur and hence, the proof is complete.



FIGURE 1: Bifurcation diagram of system (3) using data (87) with varying c, the trajectory of z versus c in the range (0, 0.6).

Theorem 9. Assume that the parameter w_6 passes through the value w_6^* , where

$$w_6^* \equiv \frac{w_5(1-m)(\bar{x}+w_3\bar{y})}{w_2+(1-m)\bar{x}+c(1-m)\bar{y}} - \frac{w_7}{w_8}.$$
 (69)

Then, system (3) near the PFEP undergoes a transcritical bifurcation but neither saddle-node bifurcation nor pitchfork bifurcation can occur provided that

$$(w_{2} + (1 - m)(c - w_{3})\bar{y})\Gamma_{1} + (w_{2}w_{3} - (1 - m)(c - w_{3})\bar{x})\Gamma_{2} + \frac{w_{7}\bar{N}_{1}^{2}}{w_{5}(1 - m)w_{8}^{2}} \neq 0,$$
(70)

where Γ_1, Γ_2 are given in the proof.

Proof. Note that, straightforward computation shows that as the parameter w_6 passes through the value w_6^* , the PFEP becomes a nonhyperbolic equilibrium point and the Jacobian matrix of system (3) at (E_2, w_6^*) can be written as

$$J_{2} = J(E_{2}, w_{6}^{*}) = \begin{pmatrix} -\bar{x} & -(1+w_{1})\bar{x} & -\frac{(1-m)\bar{x}}{\bar{N}_{1}} \\ w_{1}\bar{y} & 0 & -\frac{w_{3}(1-m)\bar{y}}{\bar{N}_{1}} \\ 0 & 0 & 0 \end{pmatrix}$$
(71)

where $\bar{N}_1 = w_2 + (1 - m)\bar{x} + c(1 - m)\bar{y}$. Hence, J_2 has three eigenvalues given by $\lambda_{21}^{[2]} = -(\bar{x}/2) - ((\sqrt{\bar{x}^2 - 4w_1(1 + w_1)\bar{x}\bar{y}})/2), \lambda_{22}^{[2]} = -(\bar{x}/2) + ((\sqrt{\bar{x}^2 - 4w_1(1 + w_1)\bar{x}\bar{y}})/2), \text{ and } \lambda_{23}^{[2]} = 0.$ Obviously, the first two eigenvalues, $\lambda_{21}^{[2]}$ and $\lambda_{22}^{[2]}$, have a negative real part.

Now, let $V_2 = (v_{21}, v_{22}, v_{23})^T$ be the eigenvector of J_2 corresponding to the eigenvalue $\lambda_{23}^{[2]} = 0$. Then, straightforward computation gives $V_2 = (\Gamma_1 v_{23}, \Gamma_2 v_{23}, v_{23})^T$, where $\Gamma_1 = ((w_3(1-m))/(w_1\bar{N}_1)) > 0, \Gamma_2 = -(((w_1+w_3)(1-m))/(w_1(1+w_1)\bar{N}_1)) < 0$, and v_{23} represents any nonzero real number.

Also, let $\Psi_2 = (\psi_{21}, \psi_{22}, \psi_{23})^T$ represent the eigenvector of J_2^T that corresponds to the eigenvalue $\lambda_{23}^{[2]} = 0$. Then, direct calculation shows that $\Psi_2 = (0, 0, \psi_{23})^T$, where ψ_{23} is any nonzero real number.

Since the partial derivative of vector field *F* with respect to the parameter w_6 is given by $(\partial F/\partial w_6) = F_{w_6} = (0, 0, -z)^T$, hence, we obtain $F_{w_6}(E_2, w_6^*) = (0, 0, 0)^T$. Therefore,

$$\Psi_2^{T} \left[F_{w_6}(E_2, w_6^*) \right] = 0.$$
(72)

Thus, system (3) at E_2 with $w_6 = w_6^*$ does not experience saddle-node bifurcation in view of the Sotomayor theorem. Now, we have

$$\Psi_2^{T} \left[DF_{w_6}(E_2, w_6^*) V_2 \right] = -v_{23} \psi_{23} \neq 0,$$
 (73)

where DF_{w_6} represents the derivative of F_{w_6} with respect to X. In addition, by substituting the value of E_2 , w_6^* , V_2 , and Ψ_2 in Equation (63), then we get

$$\Psi_{2}^{T} \left[D^{2} F(E_{2}, w_{6}^{*})(V_{2}, V_{2}) \right]$$

$$= \frac{2w_{5}(1-m)\psi_{23}v_{23}^{2}}{\bar{N}_{1}^{2}} \left[(w_{2} + (1-m)(c-w_{3})\bar{y})\Gamma + (w_{2}w_{3} - (1-m)(c-w_{3})\bar{x})\Gamma_{2} + \frac{w_{7}}{w_{5}(1-m)w_{8}^{2}} \right].$$
(74)

Obviously, $\Psi_2^T[D^2 F(E_2, w_6^*)(V_2, V_2)] \neq 0$ under the condition (70), and hence by the Sotomayor theorem, system (3) near the equilibrium point E_2 with $w_6 = w_6^*$ possesses a transcritical bifurcation, while pitchfork bifurcation cannot occur and hence, the proof is complete.

Theorem 10. Assume that conditions (29a) and (29b) hold. Let the parameter w_4 pass through the value $w_4^* = w_1 \hat{x} - ((w_3(1-m)\hat{z})/(w_2+(1-m)\hat{x})))$, then, system (3) near the DFEP undergoes a transcritical bifurcation but neither saddle-node bifurcation nor pitchfork bifurcation can occur provided that

$$\begin{pmatrix} w_1 \widehat{N}_1^2 + w_3 (1-m)^2 \widehat{z} \end{pmatrix} \omega_1 + c w_3 (1-m)^2 \widehat{z} - w_3 (1-m) \widehat{N}_1 \omega_2 \neq 0,$$
 (75)

where ω_1, ω_2 are given in the proof.



FIGURE 2: Trajectory of system (3) using the data given by Equation (87). (a) Period 3 in Int. \mathbb{R}^3_+ where c = 0.45. (b) Time series of attractor in (a). (c) Period 5 in Int. \mathbb{R}^3_+ where c = 0.47. (d) Time series of attractor in (c). (e) Chaotic attractor in Int. \mathbb{R}^3_+ where c = 0.55. (f) Time series of attractor in (e). (g) Asymptotic stable PFEP in the *xy*-plane given by $E_2 = (0.2, 0.53, 0)$ where c = 0.6.

Proof. Straightforward computation shows that as the parameter w_4 passes through the value w_4^* then, the DFEP becomes a nonhyperbolic equilibrium point and the Jacobian matrix of system (3) at (E_3, w_4^*) can be written as

$$J_{3} = J(E_{3}, w_{4}^{*}) = \left[\hat{b}_{ij}\right]_{3\times 3},$$
(76)

where $\hat{b}_{ij} = b_{ij}$ for all i, j = 1, 2, 3 are given in Equation (28) with $\hat{b}_{22}(w_4^*) = 0$. Hence, J_3 has three eigenvalues given by

$$\begin{split} \lambda_{31}^{[3]} &= -\frac{(b_{11}+b_{33})}{2} - \frac{\sqrt{(b_{11}+b_{33})^2 - 4(b_{11}b_{33}-b_{13}b_{31})}}{2}, \\ \lambda_{32}^{[3]} &= 0, \\ \lambda_{33}^{[3]} &= -\frac{(b_{11}+b_{33})}{2} + \frac{\sqrt{(b_{11}+b_{33})^2 - 4(b_{11}b_{33}-b_{13}b_{31})}}{2}. \end{split}$$

Clearly, the eigenvalues $\lambda_{31}^{[3]}$ and $\lambda_{33}^{[3]}$ have negative real parts under the conditions (29a) and (29b).

Now, let $V_3 = (v_{31}, v_{32}, v_{33})^T$ be the eigenvector of J_3 corresponding to the eigenvalue $\lambda_{32}^{[3]} = 0$. Then, straightforward computation gives $V_3 = (\omega_1 v_{32}, v_{32}, \omega_2 v_{32})^T$, where $\omega_1 = ((b_{13} b_{32} - b_{12} b_{33})/(b_{11} b_{33} - b_{13} b_{31}))$ and $\omega_2 = ((b_{12} b_{31} - b_{32} b_{11})/(b_{11} b_{33} - b_{13} b_{31}))$, while v_{32} represents any nonzero real number.

Also, let $\Psi_3 = (\psi_{31}, \psi_{32}, \psi_{33})^T$ represent the eigenvector of J_3^T that corresponds to the eigenvalue $\lambda_{32}^{[3]} = 0$. Then, direct calculation shows that $\Psi_3 = (0, \psi_{32}, 0)^T$, where ψ_{32} is any nonzero real number.

Since the partial derivative of vector field F with respect to the parameter w_4 is given by $(\partial F/\partial w_4) = F_{w_4} = (0, -y, 0)^T$, hence, we obtain $F_{w_4}(E_3, w_4^*) = (0, 0, 0)^T$. Therefore,

$$\Psi_{3}^{T} \left[F_{w_{4}}(E_{3}, w_{4}^{*}) \right] = 0.$$
(78)

Thus, system (3) at E_3 with $w_4 = w_4^*$ does not experience saddle-node bifurcation in view of the Sotomayor theorem. Now, we have

$$\Psi_{3}^{T} \left[DF_{w_{4}}(E_{3}, w_{4}^{*}) V_{3} \right] = -v_{32} \psi_{32} \neq 0$$
(79)

where DF_{w_4} represents the derivative of F_{w_4} with respect to X. In addition, by substituting the value of E_3 , w_4^* , V_3 , and Ψ_3 in Equation (63), then, we get

$$\Psi_{3}^{T} \left[D^{2} F(E_{3}, w_{4}^{*}) (V_{3}, V_{3}) \right]$$

$$= \frac{2 \psi_{32} v_{32}^{2}}{\hat{N}_{1}^{2}} \left[\left(w_{1} \hat{N}_{1}^{2} + w_{3} (1-m)^{2} \hat{z} \right) \omega_{1} + c w_{3} (1-m)^{2} \hat{z} - w_{3} (1-m) \hat{N}_{1} \omega_{2} \right].$$
(80)

Obviously, $\Psi_3^{T}[D^2F(E_3, w_4^*)(V_3, V_3)] \neq 0$ under the condition (75), and hence, by the Sotomayor theorem, system (3) near the DFEP that is given by E_3 with $w_4 = w_4^*$ possesses a transcritical bifurcation, while pitchfork bifurcation cannot occur and hence, the proof is complete.

Theorem 11. *Assume that in addition to the conditions (30a), (30b), (30c), and (30e) the following conditions are satisfied*

$$\frac{a_{12}a_{31}}{a_{32}} < a_{11} < \frac{a_{12}a_{21}}{a_{22}}.$$
(81)

Then, as the parameter w_3 passes through the value w_3^* , system (3) near the CEP undergoes a saddle-node bifurcation, but neither transcritical bifurcation nor pitchfork bifurcation can occur provided that the following conditions hold

$$w_5 - \xi_4 \neq 0, \tag{82a}$$

$$\xi_3 m_{11}^{[4]} + \xi_4 m_{21}^{[4]} + m_{31}^{[4]} \neq 0, \tag{82b}$$



FIGURE 3: Bifurcation diagram of system (3) using data (87) with varying m, the trajectory of z versus m in the range (0,0.15).

where all the symbols are clearly described in the proof, while

$$w_{3}^{*} \equiv -\frac{N_{1}^{*}}{(1-m)y^{*}} \left[\frac{a_{33}(a_{11}a_{22} - a_{12}a_{21}) + a_{13}(a_{21}a_{32} - a_{31}a_{22})}{(a_{11}a_{32} - a_{12}a_{31})} \right]$$
(83)

with a_{ij} , $\forall i, j = 1, 2, 3$ are given in Equation (31) and $N_1^* = w_2 + (1 - m)x^* + c(1 - m)y^*$.

Proof. Straightforward computation shows that the above given conditions guarantee that the parameter w_3^* is positive and the determinant (D_3) of the Jacobian matrix $J_4 = J(E_4, w_3^*) = (a_{ij}(w_3^*))_{3\times 3}$ that is given in Equation (32) is zero.

Therefore, the Jacobian matrix J_4 has the following eigenvalues

$$\begin{split} \lambda_{41}^{[4]} &= -\frac{D_1}{2} + \frac{1}{2}\sqrt{D_1^2 - 4D_2}, \\ \lambda_{42}^{[4]} &= -\frac{D_1}{2} - \frac{1}{2}\sqrt{D_1^2 - 4D_2}, \\ \lambda_{43}^{[4]} &= 0. \end{split} \tag{84}$$

It is easy to verify that $\lambda_{41}^{[4]}$ and $\lambda_{42}^{[4]}$ have negative real parts under the given conditions while the third eigenvalue is zero. So the CEP becomes a nonhyperbolic point.

Let $V_4 = (v_{41}, v_{42}, v_{43})^T$ be the eigenvector of J_4 corresponding to the zero eigenvalue $\lambda_{43}^{[4]} = 0$. Then, direct computation shows that $V_4 = (\xi_1 v_{43}, \xi_2 v_{43}, v_{43})^T$, where v_{43} represents any nonzero real number and $\xi_1 = ((a_{12}a_{23} - a_{22}a_{13})/(a_{11}a_{22} - a_{12}a_{21})) < 0$ while $\xi_2 = ((a_{21}a_{13} - a_{23}a_{11})/(a_{11}a_{22} - a_{12}a_{21})) > 0$.

Let $\Psi_4 = (\psi_{41}, \psi_{42}, \psi_{43})^T$ represent the eigenvector of J_4^T that is corresponding to the eigenvalue $\lambda_{43}^{[4]} = 0$. Then, simple calculation shows that $\Psi_4 = (\xi_3 \psi_{43}, \xi_4 \psi_{43}, \psi_{43})^T$, where ψ_{43} is any nonzero real number and $\xi_3 = ((a_{21}a_{32} - a_{22}a_{31})/(a_{32}a_{33})^T)$



FIGURE 4: Trajectory of system (3) using the data given by Equation (87). (a) Period 3 in Int. \mathbb{R}^3_+ where m = 0.024. (b) Time series of attractor in (a). (c) Chaotic attractor in Int. \mathbb{R}^3_+ where m = 0.026. (d) Time series of attractor in (c). (e) Chaotic attractor in Int. \mathbb{R}^3_+ where m = 0.066. (f) Time series of attractor in (e). (g) Asymptotic stable PFEP in the *xy*-plane given by $E_2 = (0.2, 0.53, 0)$ where m = 0.15.

 $\begin{array}{l} (a_{11}a_{22}-a_{12}a_{21}))<0 \text{ while } \xi_4 = ((a_{12}a_{31}-a_{11}a_{32})/(a_{11}a_{22}-a_{12}a_{21}))>0.\\ \text{Now, we have } (\partial F/\partial w_3) = F_{w_3} = (0,((-(1-m)yz)/N_1),((w_5(1-m)yz)/(N_1)))^T, \text{ where } N_1 = w_2 + (1-m)x + c(1-m)y; \text{ hence, we obtain } F_{w_3}(E_4,w_3^*) = (0,((-(1-m)y^*z^*)/(N_1^*)),((w_5(1-m)y^*z^*)/(N_1^*)))^T. \end{array}$

Therefore, by using condition (82a) we get

$$\Psi_{4}^{T} \left[F_{w_{3}}(E_{4}, w_{3}^{*}) \right] = \psi_{43} \left(\frac{(1-m)y^{*}z^{*}}{N_{1}^{*}} \right) (w_{5} - \xi_{4}) \neq 0.$$
(85)

Consequently, the first condition of saddle-node bifurcation in view of the Sotomayor theorem is satisfied. Now, since

$$\Psi_{4}^{T} [D^{2} F(E_{4}, w_{3}^{*}) (V_{4}, V_{4})] = (\xi_{3} m_{11}^{[4]} + \xi_{4} m_{21}^{[4]} + m_{31}^{[4]}) \psi_{43},$$
(86)

where $m_{i1}^{[4]} = m_{i1}(E_4, w_3^*)$, $\forall i = 1, 2, 3$; then, according to the condition (82b), we have $\Psi_4^T[D^2F(E_4, w_3^*)(V_4, V_4)] \neq 0$. Hence, the system (3) undergoes a saddle-node bifurcation near the coexistence equilibrium point but neither



FIGURE 5: Bifurcation diagram of system (3) using data (87) with varying w_1 . (a) The trajectory of *y* versus w_1 in the range (0.3, 0.55). (b) The trajectory of *z* versus w_1 in the range (0.3, 0.55).



FIGURE 6: Trajectory of system (3) using the data given by Equation (87). (a) Periodic attractor in the *xy*-plane where $w_1 = 0.3$. (b) Time series of attractor in (a). (c) Periodic attractor in Int. \mathbb{R}^3_+ where $w_1 = 0.4$. (d) Time series of attractor in (c). (e) Chaotic attractor in Int. \mathbb{R}^3_+ where $w_1 = 0.45$. (f) Time series of attractor in (e). (g) Asymptotic stable PFEP in the *xy*-plane given by $E_2 = (0.16, 0.52, 0)$ where $w_1 = 0.6$.



FIGURE 7: Bifurcation diagram of system (3) using data (87) with varying w_2 , the trajectory of z versus w_2 in the range (0.15, 0.21).

transcritical bifurcation nor pitchfork bifurcation can occur. This completes the proof.

5. Numerical Simulation

In this section, system (3) is simulated numerically. The objectives understand the dynamical behavior of the solution of the system as the time increasing with the parameters varied and confirm our obtained analytical results. A variety of mathematical tools is used in this investigation including the bifurcation diagrams, chaotic attractors, phase portrait, and time series. Four steps of the predictor-corrector method are used to solve system (3), and then, the obtained trajectories are drawn using MATLAB version 6. Now, the following set of hypothetical parameters are used to solve system (3) numerically.

$$w_{1} = 0.5,$$

$$m = 0.1,$$

$$w_{2} = 0.2,$$

$$c = 0.55,$$

$$w_{3} = 0.4$$

$$w_{4} = 0.1,$$

$$w_{5} = 0.8,$$

$$w_{6} = 0.1,$$

$$w_{7} = 0.4,$$

$$W_{8} = 0.8$$

(87)

In order to investigate the behavior of the system (3) as a function of *c*, bifurcation diagram is drawn in Figure 1.

According to the bifurcation diagram given in Figure 1, system (3) has a complex dynamics including chaos. Moreover, the system enters to chaos through cascades of periodic doubling and then, after a specific value ($c \ge 0.58$), the system loses the persistence, and the trajectory approaches asymptotically to PFEP in the *xy*-plane as shown in Figure 2 for typical values of *c*.

The dynamical behavior of system (3) as a function of refuge rate m in the range (0,0.15) is investigated with the help of bifurcation diagram as presented in Figure 3.

Clearly, system (3) has chaotic dynamics where in between there are small periodic regions; moreover, the system approaches asymptotically to PFEP in the *xy*-plane as *m* increases above a specific value ($m \ge 0.12$), see Figure 4 for typical values of *m*.

The dynamics of system (3) as a function of w_1 in the range $(0.3 \le w_1 \le 0.55)$ is shown in the bifurcation diagram given in Figure 5 and the typical attractors given in Figure 6.

Further investigation for the effect of varying w_2 in the range (0.15, 0.21) on the dynamical behavior of system (3) is carried out using the bifurcation diagram and phase portraits with their time series at typical values of w_2 as shown in Figures 7 and 8, respectively.

According to the above figures, system (3) has a range of complex dynamics including chaos given by (0.16, 0.21), while it approaches asymptotically to PFEP otherwise.

Other parameters have been investigated too, and the obtained results are summarized in Table 1.

6. Discussion and Conclusion

In this paper, an ecoepidemiological model consisting of the prey-predator system with infectious disease in prey has been proposed and analyzed. As a defensive property against predation, it is assumed that the prey has a refuge and the predator preys upon the available prey individuals only according to the modified Holling type-II functional response. There is a harvesting process applied on the predator species only with a nonlinear harvesting rate. The existence, uniqueness, and boundedness of the solution of the system are studied. The stability analyses around all possible equilibrium points are investigated. The persistence conditions of the system are derived analytically and shown numerically. The occurrences of local bifurcation near the nonhyperbolic equilibrium points are investigated. Finally, numerical simulations are carried out to confirm our obtained analytical results and specify the control set of parameters using biologically feasible set of parameters given in (87). It is observed that the system is very sensitive to variation in the parameter set and it has a complex dynamics for a wide range of parameters including chaos. Different types of attracting sets and bifurcation diagrams are drawn to present the behavior of the system. Finally, in the following, we summarize the obtained numerical results depending on the set of data given by Equation (87).

(1) The system (3) is sensitive for varying in the preference rate of predation as shown in the bifurcation diagram represented by Figure 1, so that it has periodic dynamics in the interior of the positive octant and then periodic doubling leading to chaos, as shown in Figure 2, for the range c < 0.58. Otherwise,



FIGURE 8: Trajectory of system (3) using the data given by Equation (87). (a) Asymptotic stable PFEP in the *xy*-plane given by $E_2 = (0.2, 0.53, 0)$ where $w_2 = 0.15$. (b) Time series of attractor in (a). (c) Chaotic attractor in Int. \mathbb{R}^3_+ where $w_2 = 0.19$. (d) Time series of attractor in (c).

The parameter	The range	The dynamical behavior of system (3)
<i>w</i> ₃	$0 < w_3 < 0.39$	Approaches to periodic dynamic in <i>xy</i> -plane
	$0.39 \le w_3 < 0.42$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.42 \le w_3 < 0.64$	Approaches to periodic attractor in $Int.\mathbb{R}^3_+$
	$0.64 \le w_3 < 1$	Approaches to periodic dynamic in <i>xz</i> -plane
w_4	$0 < w_4 < 0.1$	Approaches to PFEP in <i>xy</i> -plane
	$0.1 \le w_4 < 0.11$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.11 \le w_4 < 0.18$	Approaches to periodic attractor in $Int.\mathbb{R}^3_+$
	$0.18 \le w_4 < 1$	Approaches to periodic dynamic in <i>xz</i> -plane
<i>w</i> ₅	$0 < w_5 < 0.8$	Approaches to PFEP in <i>xy</i> -plane
	$0.8 \le w_5 < 0.81$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.81 \le w_5 < 1$	Approaches to periodic attractor in $\mathrm{Int}.\mathbb{R}^3_+$
<i>w</i> ₆	$0 < w_6 < 0.1$	Approaches to periodic attractor in $Int.\mathbb{R}^3_+$
	$0.1 \le w_6 < 0.11$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.11 \le w_6 < 1$	Approaches to PFEP in <i>xy</i> -plane
<i>w</i> ₇	$0 < w_7 < 0.34$	Approaches to periodic dynamic in <i>xz</i> -plane
	$0.34 \le w_7 < 0.4$	Approaches to periodic attractor in $Int.\mathbb{R}^3_+$
	$0.4 \le w_7 < 0.41$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.41 \le w_7 < 1$	Approaches to PFEP in <i>xy</i> -plane
<i>w</i> ₈	$0 < w_8 < 0.79$	Approaches to PFEP in <i>xy</i> -plane
	$0.79 \le w_8 < 0.81$	Approaches to chaotic attractor in $Int.\mathbb{R}^3_+$
	$0.81 \le w_8 < 1$	Approaches to periodic attractor in $Int.\mathbb{R}^3_+$

TABLE 1: The dynamical behavior of system (3) using the data (87) with varying a specific parameter.

the system loses the persistence and approach asymptotically to the PFEP in the interior of the *xy*-plane

- (2) According to the bifurcation diagram as a function of the refuge rate given by Figure 3 and the attracting sets drawn in Figure 4, system (3) has a chaotic dynamics for the range m < 0.12 in between; there is windows of periodic dynamics too. Moreover, as the refuge rate increases further, the system loses their persistence and approach asymptotically to the PFEP in the interior of the *xy*-plane
- (3) According to the bifurcation diagram as a function of the infected rate given by Figure 5 and the attracting sets drawn in Figure 6, system (3) has a periodic dynamics in the interior of the *xz*-plane for the range $w_1 < 0.37$ and then, it becomes persistent in the form of periodic dynamics and periodic doubling leading to chaos for the range $w_1 \in (0.37, 0.53)$. Otherwise, system (3) loses their persistence and approach asymptotically to the PFEP in the interior of the *xy*-plane
- (4) According to the bifurcation diagram as a function of half saturation constant given by Figure 6 and the attracting sets drawn in Figure 7, system (3) has a chaotic dynamics and in between there is periodic windows for the range $w_2 \in (0.16, 0.21)$. Otherwise, system (3) loses their persistence and approach asymptotically to the PFEP in the interior of the *xy*-plane
- (5) According to the results given in Table 1, system (3) is sensitive to varying in other parameters too and still persists in the form of chaos and periodic in the interior of the positive octant for some ranges of parameters

Data Availability

The data used to support the findings of this study are mainly from the literature similar this study, where unavailable data especially values of parameters were estimated for the purpose of verifying results of mathematical analysis of model.

Conflicts of Interest

The authors declare that there are no conflicts of interest regarding the publication of this paper.

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