

Research Article **Multiple Solutions for a Class of Fractional Schrödinger-Poisson System**

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We investigate a class of fractional Schrödinger-Poisson system via variational methods. By using symmetric mountain pass theorem, we prove the existence of multiple solutions. Moreover, by using dual fountain theorem, we prove the above system has a sequence of negative energy solutions, and the corresponding energy values tend to 0. These results extend some known results in previous papers.

1. Introduction

We consider the following system via variational methods:

$$
(-\Delta)^{\alpha} w + V(x) w + \phi w
$$

= $g(x, w) + \lambda h(x) |w|^{q-2} w, \quad x \in \mathbb{R}^3$, (SP)

$$
(-\Delta)^{\beta} \phi = w^2, \quad x \in \mathbb{R}^3
$$
,

where $\lambda > 0$, $1 < q < 2$, $\alpha, \beta \in (0, 1]$, $2\beta + 4\alpha > 3$. $(-\Delta)^{\alpha}$ and $(-\Delta)$ ^{β} represent the Laplace operator of the fractional order. If $\alpha = \beta = 1$, then the system ([SP](#page-0-0)) degenerates into the standard Schrödinger-Poisson system, which describes the interaction between the same charged particles when the magnetic effect can be ignored [\[1\]](#page-6-0). In recent years, the existence, multiplicity, and centralization of solutions for the Schrödinger-Poisson system have been deeply studied via variational methods, and a great number of works have been obtained, see, for example, [\[2](#page-6-1)[–8\]](#page-6-2). On the other hand, $(-\Delta)^{\alpha}$ is a class of nonlocal pseudodiferential operators. Since nonlocal diferential equations can better and more fully describe the physical experimental phenomena than classical local diferential operators, the study of nonlinear fractional Laplace equation has become one of the most popular research felds in nonlinear analysis. In the literature [\[9](#page-6-3)], Wei considered the following system:

$$
(-\Delta)^{\alpha} w + V(x) w + \phi w = g(x, w), \quad x \in \mathbb{R}^{3},
$$

$$
(-\Delta)^{\alpha} \phi = \gamma_{\alpha} w^{2}, \quad x \in \mathbb{R}^{3}.
$$
 (1)

By using the critical point theory, the author obtained infinitely many solutions when $\alpha \in (0, 1]$. In the literature [\[10](#page-6-4)], Teng studied a system of the form

$$
(-\Delta)^{\alpha} w + V(x) w + \phi w = \theta |w|^{q-1} w + |w|^{2_{\alpha}^{*}-2} w,
$$

$$
x \in \mathbb{R}^{3}, (2)
$$

$$
(-\Delta)^{\beta} \phi = w^{2}, x \in \mathbb{R}^{3},
$$

where $q \in (1, (3 + 2\alpha)/(3 - 2\alpha))$, $\alpha, \beta \in (0, 1)$, $2\beta + 2\alpha$ > 3. In [\[11\]](#page-6-5), Zhang, Marcos, and Squassina used perturbation approach to obtain the existence of solutions for the following system when the nonlinear term is subcritical or critical

$$
(-\Delta)^{\alpha} w + \lambda \phi w = g(w), \quad x \in \mathbb{R}^{3},
$$

$$
(-\Delta)^{\beta} \phi = \lambda w^{2}, \quad x \in \mathbb{R}^{3},
$$

$$
(3)
$$

where $\lambda > 0$, $\alpha, \beta \in [0, 1]$. In [\[12\]](#page-6-6), Duarte and Souto investigated the following system via variational methods

$$
(-\Delta)^{\alpha} w + V(x) w + \phi w = g(w), \quad x \in \mathbb{R}^{3},
$$

$$
(-\Delta)^{\beta} \phi = w^{2}, \quad x \in \mathbb{R}^{3},
$$

$$
(4)
$$

where $\alpha \in (3/4, 1), \beta \in (0, 1), V : \mathbb{R}^3 \longrightarrow \mathbb{R}$ is a periodic potential. A positive solution and a ground state solution were got in [\[12](#page-6-6)]. In [\[13\]](#page-6-7), Li studied a system of the following form:

$$
(-\Delta)^{\alpha} w + V(x) w + \phi w = g(x, w), \quad x \in \mathbb{R}^{3},
$$

$$
(-\Delta)^{\beta} \phi = w^{2}, \quad x \in \mathbb{R}^{3},
$$

$$
(5)
$$

where α , $\beta \in (0, 1]$, $2\beta + 4\alpha > 3$. Combining the perturbation method with mountain pass theorem, the existence of nontrivial solutions was obtained in [\[13](#page-6-7)]. In [\[14\]](#page-7-0), Yu, Zhao, and Zhao studied the following fractional Schrödinger–Poisson system with critical growth via variational methods

$$
\varepsilon^{2\alpha} \left(-\Delta\right)^{\alpha} w + V(x) w + \phi w = |w|^{2_{\alpha}^{*}-2} w + g(w),
$$

$$
x \in \mathbb{R}^{3}, \quad (6)
$$

$$
\varepsilon^{2\alpha} \left(-\Delta\right)^{\alpha} \phi = w^{2}, \quad x \in \mathbb{R}^{3},
$$

where $\alpha \in (3/4, 1), 2^*_{\alpha} = 6/(3 - 2\alpha)$, the potential V is continuous with positive infimum, g is continuous and subcritical at infnity. Under some Monotone hypothesis on , the existence of positive ground state solution is got in [\[14](#page-7-0)]. For small $\varepsilon > 0$, a multiple result is also got in [\[14\]](#page-7-0).

Inspired by [\[9](#page-6-3)[–16\]](#page-7-1), in this paper, we prove the existence of multiple solutions for system ([SP](#page-0-0)) by symmetric mountain pass theorem. Moreover, we prove the system ([SP](#page-0-0)) has a sequence of negative energy solutions by dual fountain theorem. The assumptions on V and nonlinearity q in this paper are given below:

(V) $V \in C(\mathbb{R}^3, \mathbb{R})$, $V_0 := \inf_{x \in \mathbb{R}^3} V(x) > 0$ and $\lim_{|x| \to +\infty} V(x) = +\infty;$

(H1) $g \in C(\mathbb{R}^3 \times \mathbb{R}, \mathbb{R})$, and there exists $C_1 > 0$ such that $|g(x, w)| \le C_1(|w| + |w|^{p-1})$, where $p \in (4, 2^*_{\alpha}), 2^*_{\alpha} =$ $6/(3 - 2\alpha);$

(H2) there exist κ > 4 and r > 0 such that 0 < $\kappa G(x, w) = \kappa \int_0^w g(x, t) dt \le g(x, w)w$, for $|w| \ge r$. Moreover, $\inf_{x \in \mathbb{R}^3, |w|=r} G(x, w) > 0;$

(H3) $g(x, -w) = -g(x, w), x \in \mathbb{R}^3, w \in \mathbb{R}$; (H4) $h : \mathbb{R}^3 \longrightarrow \mathbb{R}^+$, and $h \in L^{2/(2-q)}(\mathbb{R}^3)$.

2. Preliminaries

For $1 \leq v < \infty$, $L^{\nu}(\mathbb{R}^3)$ denotes the usual Lebesgue space with norm $||w||_{\nu} = (\int_{\mathbb{R}^3} |w|^{\nu} dx)^{1/\nu}$. Fix $\alpha \in (0, 1)$, fractional Sobolev space $H^{\alpha}(\mathbb{R}^{3})$ denoted as

$$
H^{\alpha}(\mathbb{R}^{3})
$$

 := { $\omega \in L^{2}(\mathbb{R}^{3}): \int_{\mathbb{R}^{3}} (1 + |l|^{2\alpha}) |\mathcal{F}w(l)|^{2} dl < \infty$ }, (7)

equipped with the norm

$$
\|w\|_{H^{\alpha}} = \left(\int_{\mathbb{R}^3} \left(|\mathcal{F}w\left(l\right)|^2 + |l|^{2\alpha} \left|\mathcal{F}w\left(l\right)\right|^2\right) dl\right)^{1/2},\qquad(8)
$$

where $\mathcal{F}w$ denotes the Fourier transform of function w . Let $g \in C_0^{\infty}(\mathbb{R}^3)$; the fractional Laplacian operator $(-\Delta)^{\alpha}$: $C_0^{\infty}(\mathbb{R}^3) \longrightarrow (C_0^{\infty}(\mathbb{R}^3))'$ is defined by

$$
(-\Delta)^{\alpha} g = \mathcal{F}^{-1} (|l|^{2\alpha} (\mathcal{F}g)), \quad l \in \mathbb{R}^3. \tag{9}
$$

According to Plancherel theorem [\[17\]](#page-7-2), one has $\|\mathcal{F}w\|_2$ = $\|w\|_2$, $\|I\|^{\alpha} \mathcal{F}w\|_2 = \|(-\Delta)^{\alpha/2} w\|_2$. By [\(8\),](#page-1-0) we define the equivalent norm

$$
\|w\|_{H^{\alpha}} = \left(\int_{\mathbb{R}^3} \left(\left|(-\Delta)^{\alpha/2} w(x)\right|^2 + \left|w(x)\right|^2\right) dx\right)^{1/2}.
$$
 (10)

 $D^{\alpha,2}(\mathbb{R}^3)$ is denoted as

$$
D^{\alpha,2}\left(\mathbb{R}^3\right) = \left\{w \in L^{2^*_{\alpha}}\left(\mathbb{R}^3\right) : \left|l\right|^{\alpha} \mathcal{F}w\left(l\right) \in L^2\left(\mathbb{R}^3\right)\right\}. \quad (11)
$$

In particular, $D^{\alpha,2}(\mathbb{R}^3)$ is the completion of $C_0^{\infty}(\mathbb{R}^3)$, with respect to the norm

$$
\|w\|_{D^{\alpha,2}} = \left(\int_{\mathbb{R}^3} \left|(-\Delta)^{\alpha/2} w\right|^2 dx\right)^{1/2}
$$

$$
= \left(\int_{\mathbb{R}^3} |l|^{2\alpha} |\mathcal{F}w(l)|^2 dl\right)^{1/2}.
$$
 (12)

 $L^q(\mathbb{R}^3, h)$, $q \in (1, 2)$ represents a weighted Lebesgue space, that is,

$$
L^{q}(\mathbb{R}^{3}, h) = \left\{ w : \mathbb{R}^{3} \longrightarrow \mathbb{R} \text{ is measurable and } \int_{\mathbb{R}^{3}} h(x) |w|^{q} dx \qquad (13)
$$

$$
< +\infty \left\},
$$

equipped with the norm

$$
\|w\|_{L^{q}(\mathbb{R}^{3},h)} = \left(\int_{\mathbb{R}^{3}} h(x) |w|^{q} dx\right)^{1/q}.
$$
 (14)

For convenience, we use C to represent any positive constants which may change from line to line. According to [\[18](#page-7-3)], the embedding $H^{\alpha}(\mathbb{R}^3) \hookrightarrow L^{\nu}(\mathbb{R}^3)$ is continuous for all $\nu \in [2,2^*_\alpha]$, i.e., there exists $M_\nu > 0$ satisfying

$$
\|w\|_{\nu} \le M_{\nu} \|w\|_{H^{\alpha}}, \quad w \in H^{\alpha}\left(\mathbb{R}^{3}\right). \tag{15}
$$

So, the embedding $H^{\alpha}(\mathbb{R}^3) \hookrightarrow L^{\nu}(\mathbb{R}^3)$ is continuous when $2\beta + 4\alpha > 3$. Fix $w \in H^{\alpha}(\mathbb{R}^3)$; we define the nonlinear operator $L_w : D^{\beta,2}(\mathbb{R}^3) \longrightarrow \mathbb{R}$ by

$$
L_w(v) = \int_{\mathbb{R}^3} w^2 v dx.
$$
 (16)

Hence

$$
|L_{w}(v)| \leq \left(\int_{\mathbb{R}^{3}} |w(x)|^{12/(3+2\beta)} dx\right)^{(3+2\beta)/6}
$$

$$
\cdot \left(\int_{\mathbb{R}^{3}} |v(x)|^{2_{\beta}^{\sharp}} dx\right)^{1/2_{\beta}^{\sharp}}
$$

$$
\leq C \left(\int_{\mathbb{R}^{3}} |w(x)|^{12/(3+2\beta)} dx\right)^{(3+2\beta)/6} ||v||_{D^{\beta,2}}
$$

$$
\leq C ||w||^{2} ||v||_{D^{\beta,2}}.
$$

$$
(17)
$$

By Lax-Milgram theorem, we can find a unique $\phi_w^{\beta} \in$ $D^{\beta,2}(\mathbb{R}^3)$ such that

$$
\int_{\mathbb{R}^3} \left(-\Delta\right)^{\beta/2} \phi_w^{\beta} \left(-\Delta\right)^{\beta/2} v dx = \int_{\mathbb{R}^3} w^2 v dx,
$$
\n
$$
\forall v \in D^{\beta,2} \left(\mathbb{R}^3\right),
$$
\n(18)

and ϕ_w^{β} is expressed as

$$
\phi_w^{\beta}(x) = c_{\beta} \int_{\mathbb{R}^3} \frac{w^2(y)}{|x - y|^{3 - 2\beta}} dy,
$$
\n(19)

\nwhere $c_{\beta} = \frac{\Gamma(3/2 - 2\beta)}{\pi^{3/2} 2^{2\beta} \Gamma(\beta)}$.

According to [\(19\),](#page-2-0) $\phi_w^{\beta} \ge 0$ for all $x \in \mathbb{R}^3$. Since $\beta \in (0,1]$, $2\beta + 4\alpha > 3$, we can also get $12/(3 + 2\beta) \in (2, 2^*_\alpha)$. Together with [\(17\)](#page-2-1) and [\(18\),](#page-2-2)

$$
\|\phi_w^{\beta}\|_{D^{\beta,2}} = \int_{\mathbb{R}^3} |(-\Delta)^{\beta/2} \phi_w^{\beta}|^2 dx
$$

\n
$$
\leq C \left(\int_{\mathbb{R}^3} |w(x)|^{12/(3+2\beta)} dx \right)^{(3+2\beta)/6}
$$
 (20)
\n
$$
\leq C \|w\|^2.
$$

From Hölder's inequality and [\(19\),](#page-2-0)

$$
\int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx
$$
\n
$$
\leq C \left(\int_{\mathbb{R}^3} |w(x)|^{(12/(3+2\alpha))dx} \right)^{(3+2\alpha)/6} \left\| \phi_w^{\beta} \right\|_{D^{\beta,2}} \tag{21}
$$
\n
$$
\leq C \left\| w \right\|^2 \left\| \phi_w^{\beta} \right\|_{D^{\beta,2}} \leq C \left\| w \right\|^4.
$$

Evidently,

$$
\int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx \le C ||w||^4.
$$
 (22)

Substituting [\(18\)](#page-2-2) into system ([SP](#page-0-0)), system ([SP](#page-0-0)) is equivalent to

$$
(-\Delta)^{\alpha} w + V(x) w + \phi_w^{\beta} w
$$

= $g(x, w) + \lambda h(x) |w|^{q-2} w.$ (S)

For the equation (S) (S) (S) , we define the work space E as

$$
E \coloneqq \left\{ w \in H^{\alpha}\left(\mathbb{R}^3\right) : \int_{\mathbb{R}^3} V\left(x\right) w^2 dx < \infty \right\},\tag{23}
$$

with

$$
\|w\| = \left(\int_{\mathbb{R}^3} \left(\left|(-\Delta)^{\alpha/2} w\right|^2 + V(x) w^2\right) dx\right)^{1/2}.\tag{24}
$$

Lemma 1 (see [\[19](#page-7-4)]). *Assume condition (V) holds,* $\alpha \in (0, 1)$ *,* $\nu \in [2, 2^*_{\alpha})$; then the embedding $E \hookrightarrow L^{\nu}(\mathbb{R}^3)$ is compact.

From (V) and (H1)-(H4),
$$
I: E \longrightarrow \mathbb{R}
$$

\n
$$
I(w) = \frac{1}{2} \int_{\mathbb{R}^3} (|(-\Delta)^{\alpha/2} w|^2 + V(x) w^2) dx
$$
\n
$$
+ \frac{1}{4} \int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx - \int_{\mathbb{R}^3} G(x, w) dx \qquad (25)
$$
\n
$$
- \frac{\lambda}{q} \int_{\mathbb{R}^3} h(x) |w|^q dx.
$$

is well defined. Moreover, by Lemma [1,](#page-2-4) $I \in C^1(E, \mathbb{R})$ with

$$
\langle I'(w), v \rangle
$$

= $\int_{\mathbb{R}^3} ((-\Delta)^{\alpha/2} w (-\Delta)^{\alpha/2} v + V(x) wv) dx$
+ $\int_{\mathbb{R}^3} \phi_w^{\beta} wv dx - \int_{\mathbb{R}^3} g(x, w) v dx$
- $\lambda \int_{\mathbb{R}^3} h(x) |w|^{q-2} wv dx, v \in E$ (26)

Proposition 2. *(i) If equation* ([S](#page-2-3)) *has a solution* $w \in E$ *, then system* ([SP](#page-0-0)) has a solution $(w, \phi) \in E \times D^{\beta,2}(\mathbb{R}^3)$ *.*

(ii) If for every $v \in E$ *, the following equation*

$$
\int_{\mathbb{R}^3} \left(\left(-\Delta \right)^{\alpha/2} w \left(-\Delta \right)^{\alpha/2} v + V \left(x \right) wv \right) dx
$$

+
$$
\int_{\mathbb{R}^3} \phi_w^{\beta} wv dx - \int_{\mathbb{R}^3} g \left(x, w \right) v dx \qquad (27)
$$

-
$$
\lambda \int_{\mathbb{R}^3} h \left(x \right) |w|^{q-2} wv dx = 0
$$

holds, then $w \in E$ *is a solution of* ([S](#page-2-3)).

Set ${e_i}_{i=1}^{\infty}$ as a set of normalized orthogonal basis of E, $X_i = Re_i \cdot Y_k = \bigoplus_{i=1}^k X_i, \ Z_k = \overline{\bigoplus_{i=k+1}^\infty X_i}, \ k \in \mathbb{N}$. Obviously, $E = Y_k \oplus Z_k.$

Definition 3 (see [\[20,](#page-7-5) [21](#page-7-6)]). Set $I \in C^1(E, \mathbb{R})$, $c \in \mathbb{R}$. If any sequence $\{w_n\} \subset E$ satisfying

$$
I(w_n) \longrightarrow c,
$$

\n
$$
I'(w_n) \longrightarrow 0,
$$

\n
$$
n \longrightarrow \infty,
$$

\n(28)

has a convergent subsequence, then I satisfies the $(PS)_{c}$ condition. Any sequence satisfying [\(28\)](#page-2-5) is called the (PS) _c sequence.

Definition 4 (see [\[21](#page-7-6), [22\]](#page-7-7)). Set $I \in C^1(E, \mathbb{R})$, $c \in \mathbb{R}$. If every sequence $\{w_{n_j}\} \subset E$, satisfying

$$
w_{n_j} \in Y_{n_j},
$$

\n
$$
I(w_{n_j}) \longrightarrow c,
$$

\n
$$
I|_{Y_{n_j}}' \longrightarrow 0,
$$

\n
$$
n \longrightarrow \infty,
$$

\n(29)

has a convergent subsequence, then *I* satisfies the $(PS)_c^*$ condition.

Proposition 5 (see [\[23\]](#page-7-8)). *Let* $E = Y \oplus Z$ *be a Banach space* $with \dim Y < \infty$. Assume $I \in C^1(E, \mathbb{R})$ *is an even functional and satisfies the* (PS)_{*c}* condition and</sub>

(A1) there exist ω , μ > 0 *satisfying* $I_{\partial B_{\alpha} \cap Z}$ $\inf_{w \in Z, \|w\| = \omega} I(w) \geq \mu;$

(A2) for every linear subspace $U \subset E$ *with* dim $U \subset \infty$ *, there exists a constant* $L = L(U)$ *such that* $\max_{w \in U, ||w|| \ge L} I(w)$ < 0*,*

then has a list of unbounded critical points.

Proposition 6 (see [\[22\]](#page-7-7)). Assume that $I \in C^1(E, \mathbb{R})$ is an even *functional,* $k_0 \in \mathbb{N}$. If for any $k > k_0$, there exist $r_k > \gamma_k > 0$ *such that*

(C1) $b_k = \inf\{I(w) : w \in Z_k, ||w|| = r_k\} \geq 0;$ (C2) $a_k = \max\{I(w) : w \in Y_k, ||w|| = \gamma_k\} < 0;$ (C3) $c_k = \inf\{I(w) : w \in Z_k, ||w|| \leq r_k\} \longrightarrow 0, k \longrightarrow \infty;$ (C4) *for every* $c \in [c_{k_0}, 0]$, *I satisfies the* $(PS)_c^*$ *condition*,

then has a sequence of negative critical points that converge to 0*.*

3. Main Results

Lemma 7. *If hypotheses (V) and (H1)-(H4) hold, then for any* $c \in \mathbb{R}$, *I satisfies the* $(PS)_c$ *condition.*

Proof. First, we prove the $(PS)_c$ sequence $\{w_n\}$ of I is bounded. According to (H4), it is easy to get that

$$
\int_{\mathbb{R}^3} h(x) |w_n|^q dx
$$
\n
$$
\leq \left(\int_{\mathbb{R}^3} |h(x)|^{2/(2-q)} dx \right)^{(2-q)/2} \left(\int_{\mathbb{R}^3} |w_n|^2 dx \right)^{q/2} \quad (30)
$$
\n
$$
\leq \|h\|_{2/(2-q)} \|w\|_2^q \leq M_2^q \|h\|_{2/(2-q)} \|w_n\|^q.
$$

By condition (H2), there exists $r > 0$ such that

$$
g(x, w) w \ge \kappa G(x, w), \quad |w| \ge r. \tag{31}
$$

Moreover, for any given $C_0 \in (0, (1/16)V_0)$, we can choose a constant $\delta > 0$ such that

$$
\left|\frac{1}{\kappa}g\left(x,w\right)w - G\left(x,w\right)\right| \le C_0 w^2, \text{ for } |w| \le \delta. \tag{32}
$$

From condition (H1), when $\delta \le |u| \le r$, one has

$$
\left|\frac{1}{\kappa}g\left(x,w\right)w - G\left(x,w\right)\right| \le C\left(\frac{1}{\delta} + r^{p-2}\right)w^2. \tag{33}
$$

So, for any $|w| \leq r$,

$$
\left|\frac{1}{\kappa}g\left(x,w\right)w - G\left(x,w\right)\right| \leq C_0w^2 + C\left(\frac{1}{\delta} + r^{p-2}\right)w^2. \tag{34}
$$

For $\lim_{|x| \to +\infty} V(x) = +\infty$, there exists $L > r > 0$ such that

$$
\frac{1}{16}V(x) > C\left(\frac{1}{\delta} + r^{p-2}\right), \quad |x| \ge L. \tag{35}
$$

Now [\(34\)](#page-3-0) implies

$$
\frac{1}{4} \int_{\mathbb{R}^3} V(x) w_n^2 dx
$$
\n
$$
+ \int_{|w_n(x)| \le r} \left[\frac{1}{\kappa} g(x, w_n) w_n - G(x, w_n) \right] dx
$$
\n
$$
\ge \frac{1}{4} \int_{\mathbb{R}^3} V(x) w_n^2 dx
$$
\n
$$
- \int_{|w_n(x)| \le r} \left[C_0 |w_n|^2 + C \left(\frac{1}{\delta} + r^{p-2} \right) |w_n|^2 \right] dx
$$
\n
$$
= \frac{1}{8} \int_{\mathbb{R}^3} V(x) w_n^2 dx - \int_{|w_n(x)| \le r} C_0 |w_n|^2 dx
$$
\n
$$
+ \frac{1}{8} \int_{\mathbb{R}^3} V(x) w_n^2 dx \qquad (36)
$$
\n
$$
- \int_{|w_n(x)| \le r} C \left(\frac{1}{\delta} + r^{p-2} \right) |w_n|^2 dx
$$
\n
$$
\ge \int_{|w_n(x)| \le r} \left[\frac{1}{16} V_0 - C_0 \right] w_n^2 dx
$$
\n
$$
+ \frac{1}{8} \int_{\mathbb{R}^3} V(x) w_n^2 dx - C \left(\frac{1}{\delta} + r^{p-2} \right) r^2
$$
\n
$$
\cdot \text{meas} \{x \in \mathbb{R}^3 \mid |x| \le L \} \ge \frac{1}{8} \int_{\mathbb{R}^3} V(x) w_n^2 dx
$$
\n
$$
- C \left(\frac{1}{\delta} + r^{p-2} \right) r^2 \cdot \text{meas} \{x \in \mathbb{R}^3 \mid |x| \le L \}.
$$

Since ${w_n}$ is the $(PS)_c$ sequence, when *n* is large enough,

$$
c + ||w_n|| \ge I (w_n) - \frac{1}{\kappa} \langle I'(w_n), w_n \rangle
$$

\n
$$
= \left(\frac{1}{2} - \frac{1}{\kappa}\right) \int_{\mathbb{R}^3} (|\nabla w_n|^2 + V(x) w_n^2) dx
$$

\n
$$
+ \left(\frac{1}{4} - \frac{1}{\kappa}\right) \int_{\mathbb{R}^3} \phi_{w_n}^2 w_n^2 dx
$$

\n
$$
+ \int_{\mathbb{R}^3} \left[\frac{1}{\kappa} g(x, w_n) w_n - G(x, w_n)\right] dx
$$

\n
$$
+ \left(\frac{1}{\kappa} - \frac{1}{q}\right) \lambda \int_{\mathbb{R}^3} h(x) |w_n|^q dx
$$

\n
$$
\ge \frac{1}{4} \int_{\mathbb{R}^3} (|\nabla w_n|^2 + V(x) w_n^2) dx
$$

\n
$$
+ \int_{\mathbb{R}^3} \left[\frac{1}{\kappa} g(x, w_n) w_n - G(x, w_n)\right] dx
$$

\n
$$
- \left(\frac{1}{q} - \frac{1}{\kappa}\right) \lambda M_2^q ||h||_{2/(2-q)} ||w_n||^q
$$

\n
$$
\ge \frac{1}{4} \int_{\mathbb{R}^3} |\nabla w_n|^2 dx + \frac{1}{4} \int_{\mathbb{R}^3} V(x) w_n^2 dx
$$

\n
$$
+ \int_{|w_n(x)| \le r} \left[\frac{1}{\kappa} g(x, w_n) w_n - G(x, w_n)\right] dx
$$

\n
$$
- \left(\frac{1}{q} - \frac{1}{\kappa}\right) \lambda M_2^q ||h||_{2/(2-q)} ||w_n||^q.
$$

Thus, according to (36) , when n is large enough,

$$
c + ||w_n||
$$

\n
$$
\geq \frac{1}{4} \int_{\mathbb{R}^3} |\nabla w_n|^2 dx + \frac{1}{4} \int_{\mathbb{R}^3} V(x) w_n^2 dx
$$

\n
$$
+ \int_{|w_n(x)| \leq r} \left[\frac{1}{\kappa} g(x, w_n) w_n - G(x, w_n) \right] dx
$$

\n
$$
- \left(\frac{1}{q} - \frac{1}{\kappa} \right) \lambda M_2^q ||h||_{2/(2-q)} ||w_n||^q
$$

\n
$$
\geq \frac{1}{8} ||w_n||^2 - C \left(\frac{1}{\delta} + r^{p-2} \right) r^2
$$

\n
$$
\cdot meas \{x \in \mathbb{R}^3 | |x| \leq L \}
$$

\n
$$
- \left(\frac{1}{q} - \frac{1}{\kappa} \right) \lambda M_2^q ||h||_{2/(2-q)} ||w_n||^q.
$$
 (38)

Therefore, $\{w_n\}$ is a bounded sequence in E. By Lemma [1,](#page-2-4) there exists $\widehat{w} \in E$ such that

$$
w_n \rightharpoonup \widehat{w} \quad \text{in } E,
$$

\n
$$
w_n \rightharpoonup \widehat{w} \quad \text{in } L^{\nu}(\mathbb{R}^3),
$$

\n
$$
w_n(x) \rightharpoonup \widehat{w}(x) \quad \text{a.e. on } \mathbb{R}^3.
$$

\n(39)

Next, we define the linear operator $B_{\varphi}: E \longrightarrow \mathbb{R}$ as

$$
B_{\varphi}(\nu) = \left(-\Delta\right)^{\alpha/2} \varphi \left(-\Delta\right)^{\alpha/2} \nu. \tag{40}
$$

From Hölder's inequality, we obtain

$$
\left|B_{\varphi}\left(v\right)\right| \le \left\|\varphi\right\| \left\|v\right\|, \quad v \in E. \tag{41}
$$

Now by Lemma [1](#page-2-4) and [\(22\),](#page-2-6)

$$
\|\int_{\mathbb{R}^{3}}\phi_{w_{n}}^{\beta}w_{n}(w_{n}-\widehat{w})dx\|
$$
\n
$$
\leq \|\phi_{w_{n}}^{\beta}\|_{2_{\beta}^{*}}\|w_{n}\|_{12/(3+2\beta)}\|w_{n}-\widehat{w}\|_{12/(3+2\beta)}
$$
\n
$$
\leq C\|\phi_{w_{n}}^{\beta}\|_{\mathscr{D}^{\beta,2}}\|w_{n}\|_{12/(3+2\beta)}\|w_{n}-\widehat{w}\|_{12/(3+2\beta)}
$$
\n
$$
\leq C\|w_{n}\|_{12/(3+2\beta)}^{3}\|w_{n}-\widehat{w}\|_{12/(3+2\beta)}
$$
\n
$$
\leq C\|w_{n}\|^{3}\|w_{n}-\widehat{w}\|_{12/(3+2\beta)}.
$$
\n(42)

Similarly, we can also prove

$$
\left| \int_{\mathbb{R}^3} \phi_{\widehat{w}}^{\beta} \widehat{w} \left(w_n - \widehat{w} \right) dx \right| \le C \left\| \widehat{w} \right\|^3 \left\| w_n - \widehat{w} \right\|_{12/(3+2\beta)}. \tag{43}
$$

Since $w_n \longrightarrow \widehat{w}$ in $L^{\nu}(\mathbb{R}^3)(\nu \in [2, 2^*_\alpha))$, $\lim_{n \to \infty} \int_{\mathbb{R}^3} (\phi_{w_n}^{\beta} w_n \phi_{\widehat{w}}^{\beta}(\widehat{w})(w_n - \widehat{w})dx = 0.$

At last, combining Hölder's inequality with $(H1)$ and (H4), we can easily get

$$
\lim_{n \to \infty} \int_{\mathbb{R}^3} \left(g(x, w_n) - g(x, \widehat{w}) \right) (w_n - \widehat{w}) dx = 0,
$$

$$
\lim_{n \to \infty} \int_{\mathbb{R}^3} \left(h(x) |w_n|^{q-2} w_n - h(x) |\widehat{w}|^{q-2} \widehat{w} \right)
$$

$$
\cdot \left(w_n - \widehat{w} \right) dx = 0.
$$
 (44)

Thus

$$
o(1) = \langle I'(w_n) - I'(\widehat{w}), w_n - \widehat{w} \rangle
$$

\n
$$
= B_{w_n}(w_n - \widehat{w}) - B_w(w_n - \widehat{w})
$$

\n
$$
+ \int_{\mathbb{R}^3} V(x) w_n(w_n - \widehat{w}) - V(x) \widehat{w}(w_n - \widehat{w}) dx
$$

\n
$$
+ \int_{\mathbb{R}^3} (g(x, w_n) - g(x, \widehat{w})) (w_n - \widehat{w}) dx
$$

\n
$$
+ \lambda \int_{\mathbb{R}^3} h(x) (|w_n|^{q-2} w_n - |\widehat{w}|^{q-2} w) (w_n - \widehat{w}) dx
$$

\n
$$
+ \int_{\mathbb{R}^3} (\phi_{w_n}^{\beta} w_n - \phi_{\widehat{w}}^{\beta} \widehat{w}) (w_n - \widehat{w}) dx
$$

\n
$$
= B_{w_n}(w_n - \widehat{w}) - B_w(w_n - \widehat{w})
$$

\n
$$
+ \int_{\mathbb{R}^3} V(x) w_n (w_n - \widehat{w}) - V(x) \widehat{w}(w_n - \widehat{w}) dx
$$

\n
$$
+ o(1),
$$

that is,

$$
\|w_n - \widehat{w}\|^2 = B_{w_n}(w_n - \widehat{w}) - B_w(w_n - \widehat{w})
$$

+
$$
\int_{\mathbb{R}^3} V(x) (w_n - \widehat{w})^2 dx \longrightarrow 0.
$$
 (46)

Lemma 8. *If hypotheses (V) and (H1)-(H4) hold, then satisfies* $(PS)_c^*$ *condition for all* $c \in \mathbb{R}$ *.*

Proof. By Defnition [4,](#page-3-2) we just prove the following fact: if for any $c \in \mathbb{R}, \{w_{n_j}\} \subset E$, and $w_{n_j} \in Y_{n_j}, I(w_{n_j}) \longrightarrow c, I|_{Y_{n_j}}' \longrightarrow 0$, as $n_j \longrightarrow \infty$, then $\{w_{n_j}\}$ has a convergence subsequence. The proof method is similar to Lemma [7.](#page-3-3) П

Lemma 9 (see [24]). For
$$
2 \le v < 2^*_{\alpha}
$$
, $k \in \mathbb{N}$, set
\n
$$
\beta_{\nu}(k) \coloneqq \sup \{ ||w||_{\nu} : w \in Z_k, ||w|| = 1 \},
$$
\n(47)

and then $\beta_{\nu}(k) \longrightarrow 0, k \longrightarrow \infty$ *.*

Theorem 10. *If hypotheses (V) and (H1)-(H4) hold, then we can find* $\lambda_0 > 0$, such that system (SP) has multiple solutions *for every* $\lambda < \lambda_0$. Moreover, the corresponding energy values *tend to infnity.*

Proof. According to Lemma [7,](#page-3-3) *I* satisfies (PS) _c condition. We only need to prove that I satisfies (A1) and (A2). By virtue of (H1),

$$
|G(x, w)| \le \frac{C_1}{2} |w|^2 + \frac{C_1}{p} |w|^p, \quad (x, w) \in \mathbb{R}^3 \times \mathbb{R}. \quad (48)
$$

From Lemma [9,](#page-5-0) we can get

$$
I(w) = \frac{1}{2} \int_{\mathbb{R}^3} (|(-\Delta)^{\alpha/2} w|^2 + V(x) w^2) dx + \frac{1}{4}
$$

\n
$$
\cdot \int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx - \int_{\mathbb{R}^3} G(x, w) dx - \frac{\lambda}{q}
$$

\n
$$
\cdot \int_{\mathbb{R}^3} h(x) |w|^q dx \ge \frac{1}{2} ||w||^2 - \frac{C_1}{2} \int_{\mathbb{R}^3} |w|^2 dx
$$

\n
$$
- \frac{C_1}{p} \int_{\mathbb{R}^3} |w|^p dx - M_2^q \lambda ||h||_{2/(2-q)} ||w||^q \ge \frac{1}{2} ||w||^2
$$
 (49)
\n
$$
- \frac{C_1}{2} \beta_2^2 (k) ||w||^2 - \frac{C_1}{p} M_p^p ||w||^p - M_2^q \lambda ||h||_{2/(2-q)}
$$

\n
$$
\cdot ||w||^q \ge ||w||^2 \left[\frac{1}{2} - \frac{C_1}{2} \beta_2^2 (k) - C_1 M_p^p ||w||^{p-2} - M_2^q \lambda ||h||_{2/(2-q)} ||w||^{q-2} \right].
$$

Take a sufficiently large k such that $\beta_2^2(k)$ $\langle 1/2C_1.$ Combining the above inequality, we obtain

$$
I(w) \ge ||w||^2
$$

$$
\cdot \left[\frac{1}{4} - C_1 M_p^p \, ||w||^{p-2} - M_2^q \lambda \, ||h||_{2/(2-q)} \, ||w||^{q-2} \right].
$$
 (50)

Set

$$
\eta(t) = \frac{1}{4} - C_1 M_p^p t^{p-2} - M_2^q \lambda \|h\|_{2/(2-q)} t^{q-2}, \quad t > 0. \quad (51)
$$

Since $1 < q < 2 < p$, there exists

$$
\omega_{\lambda} := \left(\frac{\lambda\left(2-q\right)M_{2}^{q} \left\|h\right\|_{2/(2-q)}}{C_{1}M_{p}^{p}\left(p-2\right)}\right)^{1/(p-q)} > 0, \qquad (52)
$$

such that $\max_{t \in \mathbb{R}^+} \eta(t) = \eta(\omega_\lambda)$. Therefore, for every $\lambda <$ $\lambda_0 := ((2-q)/4C_1M_p^p(p-q))^{(p-q)/(p-2)} \cdot (C_1(p-2)M_p^p/M_2^q(2-q))$ $q)\|h\|_{2/(2-q)}$),

$$
I(\omega) \ge \omega_{\lambda}^2 \eta(\omega_{\lambda}) = \mu > 0, \quad \text{with} \quad \|\omega\| = \omega_{\lambda}.
$$
 (53)

On the other hand, by conditions (H1) and (H2), there exist positive constants C_2, C_3 such that

$$
G(x, w) \ge C_2 |w|^{k} - C_3 |w|^{2}, \quad (x, w) \in \mathbb{R}^3 \times \mathbb{R}.
$$
 (54)

Since all the norms are equivalent in every fnite linear subspace $U \subset E$, then for $w \in U$

$$
I(w) = \frac{1}{2} \int_{\mathbb{R}^3} \left(\left| (-\Delta)^{\alpha/2} w \right|^2 + V(x) w^2 \right) dx
$$

+
$$
\frac{1}{4} \int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx - \int_{\mathbb{R}^3} G(x, w) dx
$$

-
$$
\frac{\lambda}{q} \int_{\mathbb{R}^3} h(x) |w|^q dx
$$

$$
\leq \frac{1}{2} ||w||^2 + C ||w||^4 - C_2 ||w||_K^{\kappa} - C_3 ||w||_2^2
$$

-
$$
\frac{\lambda}{q} ||w||_{L^q(\mathbb{R}^3, h)}^q.
$$
 (55)

For $q \leq 2 \leq 4 \leq \kappa$, $I(w) \longrightarrow -\infty$ as $||w|| \longrightarrow \infty$. Then there exists $L = L(U) > 0$ such that $\max_{w \in U, ||w|| \ge L} I(w) < 0$. Thus, according to Proposition [5,](#page-3-4) the system (SP) has a list of solutions $\{(w_n, \phi_n)\} \subset E \times D^{\beta,2}(\mathbb{R}^3)$, and the corresponding energy values tend to infnity. \Box

Theorem 11. *If hypotheses* (V) and (H1)-(H4) hold, then the *system (SP) has a sequence of negative energy solutions for all* $\lambda > 0$, and the energy values tend to 0.

Proof. By Lemma [8,](#page-5-1) for all $c \in \mathbb{R}$, *I* satisfies the $(PS)_{c}^{*}$ condition. It now remains to show that (Cl) - $(C3)$ are satisfied. According to Lemma [9,](#page-5-0) for every $\nu \in [2, 2^*_{\alpha}), \beta_{\nu}(k) \longrightarrow 0$, as $k \longrightarrow \infty$. Thus there exists $k_1 > 0$ such that $\beta_2(k) \le \sqrt{1/2C_1}$ for $k > k_1$. For $4 < p < 2^*_\alpha$, there exists $L \in (0, 1)$ such that

$$
\frac{1}{8} ||w||^2 \ge \frac{C_1}{p} M_p^p ||w||^p, \text{ with } ||w|| \le L. \tag{56}
$$

 $I(w) = \frac{1}{2} \int_{\mathbb{R}^3} (|(-\Delta)^{\alpha/2} w|$ $\int^2 + V(x) w^2 dx$ $+\frac{1}{4}\int_{\mathbb{R}^3} \phi_w^{\beta} w^2 dx - \int_{\mathbb{R}^3} G(x, w) dx$ $-\frac{\lambda}{q}\int_{\mathbb{R}^3} h(x)|w|^q dx$ $\geq \frac{1}{2} ||w||^2 - \frac{C_1}{2} \int_{\mathbb{R}^3} ||w||^2 dx - \frac{C_1}{p} \int_{\mathbb{R}^3} ||w||^p dx$ $-\frac{\lambda}{q}\int_{\mathbb{R}^3} h(x)\,|w|^q\,dx$ $\geq \frac{1}{2} ||w||^2 - \frac{C_1}{2} \beta_2^2(k) ||u||^2 - \frac{C_1}{p} M_p^p ||w||^p$ (57)

Hence, for $w \in Z_k$ with $||w|| \leq L$, it follows that

$$
-\frac{\lambda}{q} \int_{\mathbb{R}^3} h(x) |w|^q dx
$$

\n
$$
\geq \frac{1}{4} ||w||^2 - \frac{C_1}{p} M_p^p ||w||^p
$$

\n
$$
-\frac{\lambda}{q} ||h||_{2/(2-q)} \beta_2^q (k) ||w||^q
$$

\n
$$
\geq \frac{1}{8} ||w||^2 - \frac{\lambda}{q} ||h||_{2/(2-q)} \beta_2^q (k) ||w||^q.
$$

For every $k > k_1$, let $r_k = ((8/q)\lambda \|h\|_{2/(2-q)} \beta_2^q(k))^{1/(2-q)}$. By Lemma [9,](#page-5-0) $r_k \longrightarrow 0$, as $k \longrightarrow \infty$. Thus, there exists $k_0 > k_1$, such that for every $k \ge k_0$, $I(w) \ge 0$, for $w \in Z_k$ with $||w|| = r_k$.

Secondly, since for every fixed $k \in \mathbb{N}$, the norms are equivalent in Y_k , when k is sufficiently large, there exists a small enough γ_k such that $0 < \gamma_k < r_k$ and $I(w) < 0$ for $w \in Y_k$ with $||w|| = \gamma_k$.

Finally, according to (C3), when $k \geq k_0$, for $u \in Z_k$, with $\|w\| \leq r_k$, one has

$$
I(w) \ge -\frac{\lambda}{q} \|h\|_{2/(2-q)} \beta_2^q(k) \|w\|^q
$$

$$
\ge -\frac{\lambda}{q} \|h\|_{2/(2-q)} \beta_2^q(k) r_k^q.
$$
 (58)

Since $\beta_2(k) \longrightarrow 0$, $r_k \longrightarrow 0$, as $k \longrightarrow \infty$, therefore (C3) holds. By Proposition [6,](#page-3-5) *I* has a list of solutions $\{(w_n, \phi_n)\}\subset E\times$ $D^{\beta,2}(\mathbb{R}^3)$ such that

$$
\frac{1}{2} \int_{\mathbb{R}^3} \left(\left| (-\Delta)^{\alpha/2} w_n \right|^2 + V(x) w_n^2 \right) dx
$$

$$
+ \frac{1}{4} \int_{\mathbb{R}^3} \phi_n^{\beta} w_n^2 dx - \int_{\mathbb{R}^3} G(x, w_n) dx \qquad (59)
$$

$$
- \frac{\lambda}{q} \int_{\mathbb{R}^3} h(x) |w_n|^q dx \longrightarrow 0.
$$

Data Availability

No data were used to support this study.

Conflicts of Interest

The authors declare that they have no competing interests.

Authors' Contributions

All authors read and approved the fnal manuscript.

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References

- [1] V. Benci and D. Fortunato, "An eigenvalue problem for the Schrödinger-Maxwell equations," *Topological Methods in Nonlinear Analysis*, vol. 11, no. 2, pp. 283–293, 1998.
- [2] T. D'Aprile and D. Mugnai, "Solitary waves for nonlinear Klein-Gordon-Maxwell and Schrödinger-Maxwell equations," *Proceedings of the Royal Society of Edinburgh, Section: A Mathematics*, vol. 134, no. 5, pp. 893–906, 2004.
- [3] D. Ruiz, "The Schrödinger-Poisson equation under the effect of a nonlinear local term," *Journal of Functional Analysis*, vol. 237, no. 2, pp. 655–674, 2006.
- [4] A. Mao, L. Yang, A. Qian, and S. Luan, "Existence and concentration of solutions of Schrödinger-Possion system," *Applied Mathematics Letters*, vol. 68, pp. 8–12, 2017.
- [5] L. Zhao and F. Zhao, "Positive solutions for Schrödinger-Poisson equations with a critical exponent," *Nonlinear Analysis: Teory, Methods & Applications*, vol. 70, no. 6, pp. 2150–2164, 2009.
- [6] M. Sun, J. Su, and L. Zhao, "Solutions of a Schrödinger-Poisson system with combined nonlinearities," *Journal of Mathematical Analysis and Applications*, vol. 442, no. 2, pp. 385–403, 2016.
- [7] M. Shao and A. Mao, "Multiplicity of solutions to Schrödinger-Poisson system with concave-convex nonlinearities," *Applied Mathematics Letters*, vol. 83, pp. 212–218, 2018.
- [8] L. Wang, S. Ma, and X. Wang, "On the existence of solutions for nonhomogeneous Schrödinger-Poisson system," *Boundary Value Problems*, vol. 2016, no. 1, article no. 76, 2016.
- [9] Z. L. Wei, *Existence of Infnitely Many Solutions for the Fractional Schrodinger-Maxwell Eqautions ¨* , 2015, [https://arxiv.org/](https://arxiv.org/abs/1508.03088) [abs/1508.03088.](https://arxiv.org/abs/1508.03088)
- [10] K. Teng, "Existence of ground state solutions for the nonlinear fractional Schrödinger–Poisson system with critical Sobolev exponent," *Journal of Diferential Equations*, vol. 261, no. 6, pp. 3061–3106, 2016.
- [11] J. Zhang, J. M. do Ó, and M. Squassina, "Fractional Schrödinger-Poisson systems with a general subcritical or critical nonlinearity," *Advanced Nonlinear Studies*, vol. 16, no. 1, pp. 15–30, 2016.
- [12] R. C. Duarte and M. A. Souto, "Fractional Schrödinger-Poisson equations with general nonlinearities," *Electronic Journal of Diferential Equations*, vol. 319, pp. 1–19, 2016.
- [13] K. Li, "Existence of non-trivial solutions for nonlinear fractional Schrödinger-Poisson equations," Applied Mathematics Letters, vol. 72, pp. 1–9, 2017.

 \Box

- [14] Y. Yu, F. Zhao, and L. Zhao, "The existence and multiplicity of solutions of a fractional Schrodinger-Poisson with critical exponent," *Science China Mathematics*, vol. 61, no. 6, pp. 1039– 1062, 2018.
- [15] M. Xiang, B. Zhang, and X. Guo, "Infinitely many solutions for a fractional Kirchhoff type problem via fountain theorem," *Nonlinear Analysis. Theory, Methods & Applications. An International Multidisciplinary Journal*, vol. 120, pp. 299–313, 2015.
- [16] Q. Q. Li and X. Wu, "A new result on high energy solutions for Schrödinger-Kirchhoff type equations in \mathbb{R}^N ," *Applied Mathematics Letters*, vol. 30, pp. 24–27, 2014.
- [17] E. Di Nezza, G. Palatucci, and E. Valdinoci, "Hitchhiker's guide to the fractional Sobolev spaces," *Bulletin des Sciences Math ´ematiques*, vol. 136, no. 5, pp. 521–573, 2012.
- [18] P. Pucci, M. Xiang, and B. Zhang, "Multiple solutions for nonhomogeneous Schrödinger-Kirchhoff type equations involving the fractional *p*-Laplacian in R ," *Calculus of Variations and Partial Diferential Equations*, vol. 54, no. 3, pp. 2785–2806, 2015.
- [19] K. Teng, "Multiple solutions for a class of fractional Schrodinger equations in \mathbb{R}^N ," *Nonlinear Analysis: Real World Applications*, vol. 21, pp. 76–86, 2015.
- [20] H. Brézis, J. Coron, and L. Nirenberg, "Free vibrations for a nonlinear wave equation and a theorem of P. Rabinowitz," *Communications on Pure and Applied Mathematics*, vol. 33, no. 5, pp. 667–684, 1980.
- [21] M. Willem, *Minimax Theorems*, Birkhäuser, Boston, Mass, USA, 1996.
- [22] T. Bartsch and M. Willem, "On an elliptic equation with concave and convex nonlinearities," *Proceedings of the American Mathematical Society*, vol. 123, no. 11, pp. 3555–3561, 1995.
- [23] P. Bartolo, V. Benci, and D. Fortunato, "Abstract critical point theorems and applications to some nonlinear problems with 'strong' resonance at infinity," *Nonlinear Analysis. Theory, Methods & Applications*, vol. 7, no. 9, pp. 981–1012, 1983.
- [24] Z. Binlin, G. Molica Bisci, and R. Servadei, "Superlinear nonlocal fractional problems with infnitely many solutions," *Nonlinearity*, vol. 28, no. 7, pp. 2247–2264, 2015.

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