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Research Article

K Shell X-Ray Fluorescence Parameters of Some Elements in the Atomic Range $40 \le Z \le 50$

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The production cross sections σ_{K_i} , the fluorescence yield W_K , K shell X-ray intensity ratio I_{K_β}/I_{K_α} , the vacancy transfer probabilities from K to L shell η_{KL} , and the Γ_K level widths for some elements in the atomic range $40 \le Z \le 50$ were measured. The samples were excited by 80.998 keV gamma rays from a 10 mCi 133 Ba radioactive source. The K X-rays emitted by samples were detected by using a CdTe detector. These parameters have been theoretically calculated, also. The experimental values were compared with the theoretical and semiempirical values. Our experimental values closely agreed with theoretical values and other experimental values.

1. Introduction

Accurate experimental values of X-ray fluorescence parameters such as the production cross sections, the fluorescence yields, the intensity ratios, the vacancy transfer probabilities, and the level widths for different elements are important because of their wide use in atomic, molecular, and radiation physics, material science, environmental science, agriculture, forensic science, dosimetric computations for health physics, cancer therapy, elemental analysis, and basic studies of nuclear physics.

A vacancy in the inner shell of an atom is produced by various methods; photoionization is one of them. In this method, the incident gamma photon ejects the element electron to the continuum state, creating a vacancy in the inner shell. The fluorescence yield of an atomic shell or subshell is defined as the probability that a vacancy in that shell or subshell is filled through a radiative transition or a nonradiative transition.

In recent years, several attempts have been made for measuring X-ray fluorescence cross sections and yields for different elements. Krause et al. [1] have calculated theoretical K and L XRF cross sections. Puri et al. [2] published an extensive table of K shell X-ray fluorescence cross sections for some elements in the energy range 1–200 keV. Bennal et al.

[3] measured K X-ray fluorescence parameters for Ag, Cd, In, and Sn elements. Seven [4] has measured photon induced K X-ray cross sections for some heavy elements. Hubbell et al. [5] have collected more recent experimental values of $\sigma_{\rm K_i}$. Theoretical values of $\sigma_{\rm K_i}$ were obtained in the region 44 $\leq Z \leq$ 54 by McGuire [6, 7]. Chen et al. [8] used a Dirac-Hartree-Slater approach to calculate the σ_{K} values of elements in the atomic range $18 \le Z \le 96$. Bambynek et al. [9] have fitted their collection of selected most reliable experimental values of σ_{K_1} in the 13 $\leq Z \leq$ 92 range. The K_{α} , K_{β} , and total K X-rays fluorescence cross sections and the average fluorescence yields for six elements with $16 \le Z \le 23$ at 5.96 keV have been measured by Sahin et al. [10]. Horakeri et al. [11, 12] determined K shell fluorescence yields by using a simple method for some elements in 62 \leq Z \leq 83 at 123.6 and 320 keV energies. Pious et al. [13] determined total K fluorescence yields by using Xe filled proportional counter for some low elements at 59.5 keV. Balakrishna et al. [14] measured K fluorescence yields using HPGe low energy photon detector for some rare earth and heavy elements with 59.5 and 279.2 keV gamma rays.

Rao et al. [15] showed that the K_{β}/K_{α} intensity ratios depend on the excitation modes in 3D elements, but they could not find such dependence for the high Z elements. Dhal and Padhi [16] have investigated relative K X-ray intensities

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on the elements from Mn to Sb using 59.5 keV gamma rays. Rebohle et al. [17] have measured K_{β}/K_{α} intensity ratio for pure 3D elements. Similarly, Ertugrul et al. [18] have measured K_{β}/K_{α} intensity ratios in element range $22 \le Z \le$ 69 at 59.5 keV. Pawlowski et al. [19] have reported the valence electronic structure of Ti, Cr, Fe, and Co in some alloys from K_{β}/K_{α} X-ray intensity ratios. The vacancy transfer probabilities from K to L shell $\eta_{\rm KL}$ were obtained by measuring the K_{β}/K_{α} intensity ratios in 25 elements over the range 57 \leq $Z \le 92$ using a 25 mCi ⁵⁷Co filtered source for excitation and a Si (Li) detector by Ertugral et al. [20]. Santra et al. [21] have measured K to L shell vacancy transfer probabilities of Mo, Pd, and Cd by exciting them with the K X-rays of Ni and Sn induced by the bremsstrahlung emanating from an X-ray tube. Demir and Şahin [22] have measured K shell X-ray production cross sections and fluorescence yields for Nd, Eu, Gd, Dy, and Ho using radioisotope in the external magnetic field.

The aim of this study is to determine the production cross sections σ_{K_i} , the fluorescence yield w_K , K shell X-ray intensity ratio I_{K_β}/I_{K_α} , the vacancy transfer probabilities from K to L shell η_{KL} , and the Γ_K level widths for some elements in the atomic range $40 \le Z \le 50$. The samples were excited by 80.998 keV gamma rays from a 10 mCi 133 Ba radioactive source. The K X-rays emitted by samples were detected by using a CdTe detector.

2. Experimental Procedure

The experimental setup and the geometry used in the measurements are shown in Figure 1.

All samples were irradiated with 80.998 keV gamma rays from a 10 mCi 133 Ba radioactive source. The samples were placed at a 45° angle with respect to the direct beam, and fluorescent X-rays emitted at 90° to the direct beam were detected by the collimated CdTe detector. An AMPTEK XR-100T-CdTe detector of crystal dimensions $3\times3\times1~\text{mm}^3$ was used in this work. A 250 μm Be window is in front of the detector. The characteristic properties of CdTe detector are given in Table 1.

The preamp signals were processed by Ampteks digital processors, including PX4, DP4, and X123, using trapezoidal shaping and no rise-time discrimination. The digital processors include a multichannel analyzer, power supplies, closed-loop temperature regulation, and the bias voltage for the detector. The data were taken at 224 K and at 502 V bias. The data were acquired by the ADMCA.EXE software provided with Amptek's digital processors. Spectrum analysis was carried out using the XRFFP software, developed jointly by Amptek and by Cross Roads Scientific. Special analysis algorithms have been developed for CdTe, which are incorporated with the commercially available package. The background spectra were measured in the laboratory, using a lead filter with a thickness of 30 mm. It was observed that the background level in the lead filter is smaller than that in a room. Thus, the detector was shielded by a filter of lead to obtain a thin beam of photons and to prevent undesirable radiation. The output from the preamplifier, with pulse pileup rejection capability, was fed to a multichannel analyzer

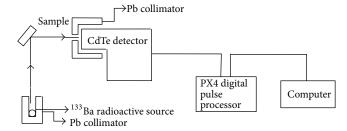


FIGURE 1: The experimental setup.

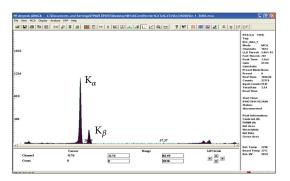


FIGURE 2: Typical K X-rays spectrum of Sn target.

interfaced with a personal computer provided with suitable software for data acquisition and peak analysis. The mean of ten channels at each side of the peaks was used to calculate the background and to define the net peak area. The peak areas were determined after the peaks were fitted either with pure Gaussians or Gaussians with exponential tails using the Microcal Origin 7.5 Software. A typical K shell X-ray spectrum of Sn target is shown in Figure 2.

The experimental production cross sections σ_{K_i} were evaluated using the relation

$$\sigma_{K_{i}} = \frac{N_{K_{i}}}{I_{0}G\varepsilon_{K}\beta_{K}m} \quad (i = \alpha, \beta), \tag{1}$$

where N_{K_i} is the measured intensity, I_0 is intensity of incident radiation, G is a geometrical factor, ε_{K_i} is the efficiency of the detector at the average K X-ray energy of the element, m is the mass per unit area of the element in g/cm^2 , and β_{K_i} is the selfabsorption correction for the target material, which accounts for the absorption in the sample of the incident photons and the emitted characteristic X-rays. The values of β_{K_i} have been calculated by using the following expression:

$$\beta_{K_i} = \frac{1 - \exp\left[-\left(\mu_1/\sin\theta + \mu_2/\sin\phi\right)m\right]}{\left(\mu_1/\sin\theta + \mu_2/\sin\phi\right)m},$$
 (2)

where μ_1 and μ_2 are the mass attenuation coefficients (from [26]) of incident photons and emitted characteristic X-rays, respectively; the angles of the incident photons and emitted X-rays with respect to the surface of samples are θ and ϕ , respectively.

The value of the factors $I_0G\varepsilon_{K_i}$, which contain terms related to the incident photon flux, geometrical factor, and the

Average electron-hole pair Mobility Average lifetime, sec Average electron-hole Atomic Density, g/cm³ Band gap, eV at 25°C, cm²/(V·sec), eV number pair eV n^* p^* 10^{-6} 6.06 1.47 48.5 4.43 1100 100

TABLE 1: The characteristic properties of CdTe detector.

efficiency of the detector, was determined by collecting the K X-ray spectra of thin samples of Zr, Nb, Mo, Ru, Rh, Pd, Ag, Cd, and Sn with the mass thickness 0.031–0.865 g/cm² in the same geometry and using the following equation:

$$I_0 G \varepsilon_{K_i} = \frac{N_{K_i}}{\sigma_{K_i} m \beta_{K_i}}, \tag{3}$$

where $N_{\rm K_i}$ is the net number of counts under the corresponding photopeak, m_i is the thickness of target in g/cm², and $\beta_{\rm K_i}$ is the self-absorption correction factor for the target material, which accounts for the absorption in the target of the incident photons and the emitted characteristic X-rays. The theoretical production cross sections $\sigma_{\rm K_i}$ were evaluated using the relation

$$\sigma_{K_i} = \sigma_K \omega_K f_{K_i}, \tag{4}$$

where $\sigma_{\rm K}$ is the K shell photoionization cross section for the given elements at the excitation energy E. The values of $\sigma_{\rm K}$ used in these calculations were taken from the report by Scofield [27]. $w_{\rm K}$ is the K shell fluorescence yield. The values of $w_{\rm K}$ were taken from the annotated bibliography of Krause [23]. $f_{\rm K_i}$ is the fractional X-ray emission rate for K_i X-rays taken from Scofield [24].

The semiempirical K shell fluorescence yields $w_{\rm K}$ were measured according to the following equation:

$$\omega_{K} = \frac{\sigma_{K}}{\sigma_{K}(E)},\tag{5}$$

where σ_K is the total K X-ray fluorescence cross section obtained experimentally and $\sigma_K(E)$ is the theoretical K shell photoionization cross section of a given element for the excitation energy E [27].

The experimental K shell X-ray intensity ratios $I_{K_{\beta}}/I_{K_{\alpha}}$ were evaluated using the equation

$$\frac{I_{K_{\beta}}}{I_{K_{\alpha}}} = \frac{N(K_{\beta})}{N(K_{\alpha})} \frac{\beta(K_{\alpha})}{\beta(K_{\beta})} \frac{\varepsilon(K_{\alpha})}{\varepsilon(K_{\beta})}, \tag{6}$$

where $N(K_{\alpha})$ and $N(K_{\beta})$ are net counts observed under the peaks corresponding to K_{α} and K_{β} X-rays, respectively, $\beta(K_{\alpha})$ and $\beta(K_{\beta})$ are the sample self-absorption correction factors for the incident photon and emitted K X-rays photons, and $\varepsilon(K_{\alpha})$ and $\varepsilon(K_{\beta})$ are photopeak efficiency values of the detector.

The experimental K to L shell total vacancy transfer probabilities $\eta_{\rm KL}$ were calculated by using the following equation [28]:

$$\eta_{\rm KL} = \frac{2 - \omega_{\rm K}}{1 + \left(I_{\rm K_{\beta}}/I_{\rm K_{\alpha}}\right)},\tag{7}$$

Table 2: The experimental $\sigma_{K_{\alpha}}$ and $\sigma_{K_{\beta}}$ production cross sections.

Sample	Experimental	Theoretical	Experimental	Theoretical
Sample	$\sigma_{ ext{K}_{lpha}}$	$\sigma_{\mathrm{K}_{lpha}}$	$\sigma_{{ m K}_eta}$	$\sigma_{{ m K}_eta}$
Zr	118.374	118.374	20.561	20.562
Nb	137.257	137.257	24.322	24.322
Mo	148.251	148.252	26.818	26.819
Ru	172.576	172.577	32.358	32.358
Rh	202.336	202.336	38.565	38.565
Pd	222.744	222.745	43.056	43.057
Ag	244.202	244.203	47.961	47.961
Cd	267.368	267.369	55.016	53.340
Sn	316.684	316.685	64.274	65.269

where $\omega_{\rm K}$ is the fluorescence yield and $I_{{\rm K}_\beta}/I_{{\rm K}_\alpha}$ is the intensity ratio of the K X-rays. The K level width was measured using the following relation:

$$\Gamma_{\rm K} = \frac{\Gamma_{\rm K}(R)}{w_{\rm K}},\tag{8}$$

where $\Gamma_{K}(R)$ is the radiative transition rate.

3. Results and Discussion

The overall error in the present measurements is estimated to be 4-5%. This error is the quadrature sum of the uncertainties in the different parameters used to evaluate the K shell fluorescence parameters, that is, target thickness (\leq 1%), the evaluation of the peak area (\leq 3%), detector efficiency $I_0G\varepsilon_{K_{ij}}$ (\leq 3%), and absorption correction factor (\leq 2%).

The experimental values of $\sigma_{\rm K_{\alpha}}$ and $\sigma_{\rm K_{\beta}}$ production cross sections in the atomic range $40 \le Z \le 50$ for 80.998 keV are listed in Table 2. The $\sigma_{\rm K_{\alpha}}$ and $\sigma_{\rm K_{\beta}}$ production cross sections versus atomic number were plotted in Figures 3 and 4, also.

These theoretical values were calculated by using K X-ray emission rates based on relativistic Hartree-Slater theory [24]. The present experimental results of K shell X-ray production cross sections are only compared with the calculated theoretical values. The deviations between the theoretical and experimental values are found to be 0.0006% and 3% for $\sigma_{\rm K_{\alpha}}$ and $\sigma_{\rm K_{\beta}}$, respectively. The reason for larger deviation between the experimental and the theoretical values for $\sigma_{\rm K_{\alpha}}$ is statistically poor due to lower intensity of the K $_{\beta}$ line in comparison with the K $_{\alpha}$ line.

The K shell fluorescence yields ω_K for these elements were calculated and listed in Table 3 together with theoretical values. The theoretical values were taken from Krause [23].

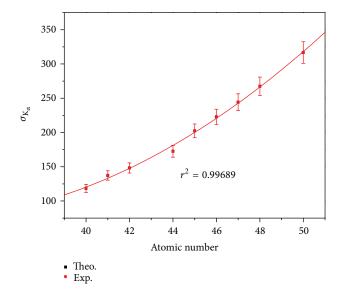


Figure 3: The variation in σ_{K_α} production cross sections as a function of atomic number.

TABLE 3: K shell X-ray fluorescence yields.

Sample	w_{K} experimental	Krause (1979) [23]
Zr	0.730	0.730
Nb	0.747	0.747
Mo	0.765	0.765
Ru	0.749	0.749
Rh	0.808	0.808
Pd	0.820	0.820
Ag	0.831	0.831
Cd	0.847	0.843
Sn	0.859	0.862

The average K shell fluorescence yields are plotted as function of the atomic number in Figure 5, also. Our experimental values have differences of 0.35–0.47% with Krause [23]. This may be due to the fact that when the elements of K X-ray energies increase, the detector efficiency decreases.

The experimental values of $I_{\rm K_{\beta}}/I_{\rm K_{\alpha}}$ X-ray intensity ratios for these elements were calculated from (6) and listed in Table 4. The experimental values of the $I_{\rm K_{\beta}}/I_{\rm K_{\alpha}}$ intensity ratios are plotted as a function of the atomic number in Figure 6. As seen from Figure 6, the $I_{\rm K_{\beta}}/I_{\rm K_{\alpha}}$ X-ray intensity ratios increase evidently with increase in atomic number. The experimental $I_{\rm K_{\beta}}/I_{\rm K_{\alpha}}$ intensity ratios are also in agreement with the theoretical values calculated for these elements as given in Table 4. Our experimental values have differences of 0.5–3% with Scofield's [24].

The experimental values of η_{KL} were calculated from (7), listed in Table 5, and plotted as a function of the atomic number in Figure 7. The values of the η_{KL} decrease with the increase in atomic numbers. K shell electrons of Zr, Nb, Mo, Ru, Rh, Pd, Ag, Cd, and Sn were excited by using the radioisotope source since K-edges energie of these elements

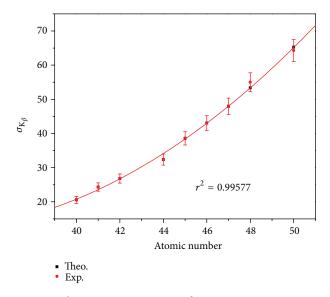


Figure 4: The variation in σ_{K_β} production cross sections as a function of atomic number.

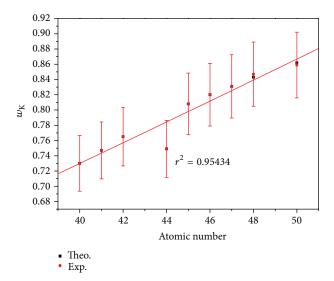


FIGURE 5: K X-ray fluorescence yields versus atomic number.

are lower than the energy of the incident photons. The vacancy created in the K shell is filled through the radiative transition or a nonradiative transition. Since the L X-ray intensity is enhanced with the K to L shell vacancy transfer, the L X-ray production cross sections of Zr, Nb, Mo, Ru, Rh, Pd, Ag, Cd, and Sn increase. As seen from Table 5, the experimental values of $\eta_{\rm KL}$ agree with the theoretically calculated values within 0.3–5%

The experimental values of $\Gamma_{\rm K}$ were calculated from (8), listed in Table 6 and plotted as a function of the atomic number in Figure 8. Our experimental values have differences of 0.2–0.5% with [25].

As a result, more experimental and theoretical values for the elements from Zr to Sn are needed for full knowledge of K X-ray fluorescence parameters. These experimental values

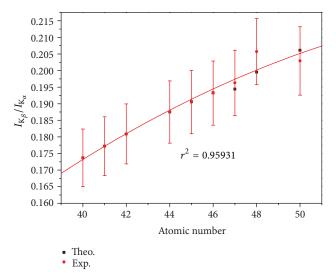


Figure 6: The variation in $I_{{\rm K}_\beta}/I_{{\rm K}_\alpha}$ as a function of atomic number.

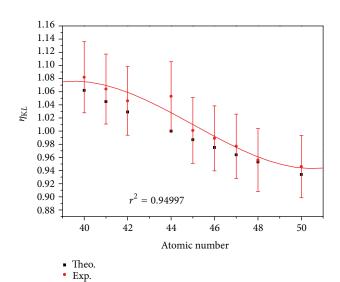


Figure 7: The variation in $\eta_{\rm KL}$ as a function of atomic number.

Table 4: $I_{{\rm K}_\beta}/I_{{\rm K}_\alpha}$ X-ray intensity ratios.

Sample	$I_{{ m K}_{eta}}/I_{{ m K}_{lpha}}$ experimental	Scofield (1974) [24]
Zr	0.1737	0.1737
Nb	0.1772	0.1772
Mo	0.1809	0.1809
Ru	0.1875	0.1875
Rh	0.1905	0.1906
Pd	0.1932	0.1933
Ag	0.1963	0.1944
Cd	0.2057	0.1995
Sn	0.2029	0.2061

are very important for qualitative and quantitative element analysis used in the XRF technique. It is well known that characteristic X-ray lines are influenced by the chemical state

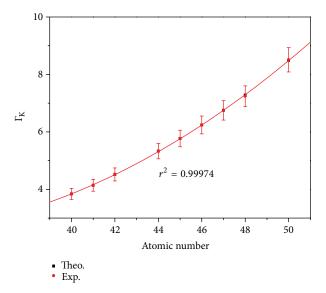


Figure 8: The variation in $\Gamma_{\!K}$ as a function of atomic number.

Table 5: K-L total vacancy transfer.

Sample	$\eta_{ ext{KL}}$ experimental	Theoretical
Zr	1.082	1.062
Nb	1.064	1.045
Mo	1.046	1.029
Ru	1.053	1.000
Rh	1.001	0.987
Pd	0.989	0.975
Ag	0.977	0.964
Cd	0.956	0.953
Sn	0.946	0.934

Table 6: K shell Γ_K level widths.

Sample	$\Gamma_{\!K}$ experimental	Krause and Oliver (1979) [25]
Zr	3.84	3.84
Nb	4.14	4.14
Mo	4.52	4.52
Ru	5.33	5.33
Rh	5.77	5.77
Pd	6.24	6.24
Ag	6.75	6.75
Cd	7.24	7.28
Sn	8.51	8.49

of the emitting element: peak energy, peak width, relative emission rates, and satellite peak formation. Also, they are affected by the detected element speciation [29]. Thus, the obtained data can be helpful for the radioisotope XRF method for elemental analysis.

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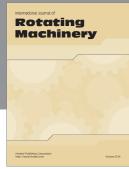
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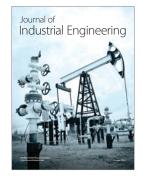
















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